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Threshold Pion Production at Large Momentum Transfers

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Abstract

We consider threshold pion electroproduction on a proton target for photon virtualities in the region 1 – 10 GeV². The S-wave multipoles at threshold, E_{0+} and L_{0+} , are calculated using light-cone sum rules.

Key words: threshold pion production, chiral symmetry

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Pion electroproduction at threshold from a proton target, $e(l) + p(P) \rightarrow e(l') + \pi^+(k) + n(P')$ and $e(l) + p(P) \rightarrow e(l') + \pi^0(k) + p(P')$, can be described in terms of two generalised form factors [1,2]

$$\begin{aligned} \langle N(P') \pi(k) | j_\mu^{em}(0) | p(P) \rangle = & \quad (1) \\ = -\frac{i}{f_\pi} \bar{N}(P') \gamma_5 \left\{ (\gamma_\mu q^2 - q_\mu \not{q}) \frac{1}{m^2} G_1^{\pi N}(Q^2) - \frac{i \sigma_{\mu\nu} q^\nu}{2m} G_2^{\pi N}(Q^2) \right\} N(P), \end{aligned}$$

which can be related to the S-wave transverse E_{0+} and longitudinal L_{0+} multipoles (at threshold). Here and below $m = 939$ MeV is the nucleon mass.

The celebrated low-energy theorem (LET) [3,4,5] relates the S-wave multipoles or, equivalently, the form factors G_1, G_2 at threshold, to the nucleon electromagnetic and axial form factors in the chiral limit $m_\pi = 0$

$$\begin{aligned} \frac{Q^2}{m^2} G_1^{\pi^0 p} &= \frac{g_A}{2} \frac{Q^2}{(Q^2 + 2m^2)} G_M^p, & G_2^{\pi^0 p} &= \frac{2g_A m^2}{(Q^2 + 2m^2)} G_E^p, & (2) \\ \frac{Q^2}{m^2} G_1^{\pi^+ n} &= \frac{g_A}{\sqrt{2}} \frac{Q^2}{(Q^2 + 2m^2)} G_M^n + \frac{1}{\sqrt{2}} G_A, & G_2^{\pi^+ n} &= \frac{2\sqrt{2} g_A m^2}{(Q^2 + 2m^2)} G_E^n. \end{aligned}$$

Here the terms in $G_{M,E}$ are due to pion emission off the initial proton state, whereas for charged pion in addition there is a contribution corresponding to the chiral rotation of the electromagnetic current [3].

The subsequent discussion concentrated mainly on the corrections to (2) due to finite pion mass [6,7]. More recently, the threshold pion production for small Q^2 was reconsidered and the low-energy theorems re-derived in the framework of the chiral perturbation theory (CHPT), see [8] for a review. The new insight gained from CHPT calculations [9] is that the expansion at small Q^2 has to be done with care as the limits $m_\pi \rightarrow 0$ and $Q^2 \rightarrow 0$ do not commute, in general. The LET predictions seem to be in good agreement with experimental data on pion photoproduction [10]. However, it appears [11,8] that the S-wave electroproduction cross section for already $Q^2 \sim 0.1 \text{ GeV}^2$ cannot be explained without taking into account chiral loops.

For larger momentum transfers the situation is much less studied as the power counting of CHPT cannot be applied. The traditional derivation of LET using PCAC and current algebra does not seem to be affected as long as the emitted pion is 'soft' with respect to the initial and final state nucleons simultaneously. The corresponding condition is, parametrically, $Q^2 \ll \Lambda^3/m_\pi$ (see, e.g. [6]) where Λ is some hadronic scale, and might be satisfied for $Q^2 \sim 1 \text{ GeV}^2$ or even higher. We are not aware of any dedicated analysis of the threshold production in the $Q^2 \sim 1 \text{ GeV}^2$ region, however.

It was suggested [12] that in the limit of very large momentum transfers the standard pQCD collinear factorisation approach becomes applicable and the helicity-conserving $G_1^{\pi N}$ form factor can be calculated for $m_\pi = 0$ in terms of chirally rotated nucleon distribution amplitudes. In practice one expects that the onset of the pQCD regime is postponed to very large momentum transfers because the factorisable contribution involves a small factor $\alpha_s^2(Q)/\pi^2$ and has to win over nonperturbative "soft" contributions that are suppressed by an extra power of Q^2 but do not involve small coefficients.

The purpose of this study is to suggest a realistic QCD-motivated model for the Q^2 dependence of the $G_{1,2}$ form factors alias S-wave multipoles at threshold in the region $Q^2 \sim 1 - 10 \text{ GeV}^2$ that can be accessible in current and future experiments in Jefferson Laboratory and elsewhere (HERMES, MAMI). In Ref. [13] we have developed a technique to calculate baryon form factors for moderately large Q^2 using light-cone sum rules (LCSR). This approach is attractive because in LCSRs "soft" contributions to the form factors are calculated in terms of the same nucleon distribution amplitudes (DAs) that enter the pQCD calculation and there is no double counting. Thus, the LCSRs provide one with the most direct relation of the hadron form factors and distribution amplitudes that is available at present, with no other nonperturbative parameters. The same technique can be applied to pion electroproduction, taking into account the semi-disconnected pion-nucleon contributions in the intermediate state. In Refs. [1,2] the G_1 and G_2 form factors are estimated in the LCSR approach for the range of momentum transfers $Q^2 \sim 1 - 10 \text{ GeV}^2$. We demonstrate that the LET results in (2) are indeed reproduced at $Q^2 \sim 1 \text{ GeV}^2$ to the required accuracy $\mathcal{O}(m_\pi)$, whereas the pQCD contribution considered in [12] formally corresponds to the leading (at large Q^2) part of the NNLO radiative correction $\sim \mathcal{O}(\alpha_s^2)$ to the sum rules. Hence our approach describes both high- Q^2 and low- Q^2 limits correctly and presents an extrapolation in between that makes maximal use of quark-hadron duality and dispersion relations.

Accurate quantitative predictions are difficult for several reasons, e.g, because the nucleon distribution amplitudes are poorly known. In order to minimise the dependence of

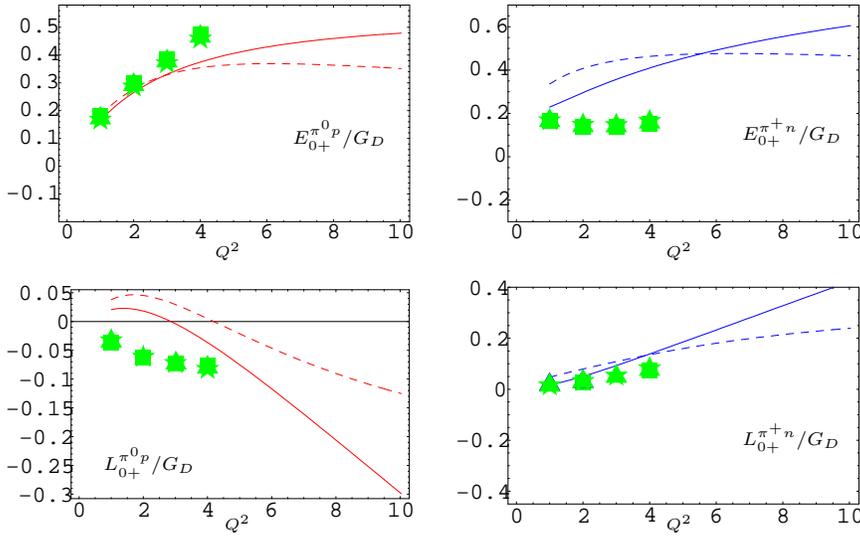


Fig. 1. The tree-level LCSR results (curves) compared to MAID07 [14] (points) for the Q^2 dependence of the electric and longitudinal partial waves at threshold, E_{0+} and L_{0+} (in units of GeV^{-1}), normalised to the dipole formula.

various parameters one may use the LCSRs to predict certain form factor ratios only, and then normalise to the electromagnetic nucleon form factors as measured in the experiment, see [2] for the details. The results are shown by the solid curves in Fig. 1, where the four partial waves at threshold are plotted as a function of Q^2 , normalised to the dipole formula $G_D(Q^2) = 1/(1 + Q^2/\mu_0^2)^2$ where $\mu_0^2 = 0.71 \text{ GeV}^2$. To give a rough idea about possible uncertainties, the “pure” LCSR predictions (all form factors and other input taken from the sum rules) are shown by dashed curves for comparison. The accuracy can be improved in future and requires calculation of radiative corrections to the LCSRs, especially if sufficiently precise lattice calculations of the moments of nucleon distribution amplitudes become available.

References

- [1] V. M. Braun, D. Yu. Ivanov, A. Lenz and A. Peters, Phys. Rev. D **75**, 014021 (2007).
- [2] V. M. Braun, D. Yu. Ivanov and A. Peters, Phys. Rev. D **77**, 034016 (2008).
- [3] N. M. Kroll and M. A. Ruderman, Phys. Rev. **93**, 233 (1954).
- [4] Y. Nambu and D. Lurie, Phys. Rev. **125**, 1429 (1962).
- [5] Y. Nambu and E. Shrauner, Phys. Rev. **128**, 862 (1962).
- [6] A. I. Vainshtein and V. I. Zakharov, Nucl. Phys. B **36**, 589 (1972).
- [7] S. Scherer and J. H. Koch, Nucl. Phys. A **534**, 461 (1991).
- [8] V. Bernard, N. Kaiser and U. G. Meissner, Int. J. Mod. Phys. E **4**, 193 (1995).
- [9] V. Bernard, N. Kaiser and U. G. Meissner, Phys. Rev. Lett. **69**, 1877 (1992).
- [10] D. Drechsel and L. Tiator, J. Phys. G **18**, 449 (1992).
- [11] V. Bernard, N. Kaiser, T. S. H. Lee and U. G. Meissner, Phys. Rev. Lett. **70**, 387 (1993).
- [12] P. V. Pobylitsa, M. V. Polyakov and M. Strikman, Phys. Rev. Lett. **87** (2001) 022001.
- [13] V. M. Braun, A. Lenz, N. Mahnke and E. Stein, Phys. Rev. D **65**, 074011 (2002); V. M. Braun, A. Lenz and M. Wittmann, Phys. Rev. D **73**, 094019 (2006).

[14] D. Drechsel, S. S. Kamalov and L. Tiator, Eur. Phys. J. A **34**, 69 (2007).