Durham Research Online

Deposited in DRO:
06 October 2016

Version of attached file:
Published Version

Peer-review status of attached file:
Peer-reviewed

Citation for published item:

Further information on publisher’s website:
http://dx.doi.org/10.1093/mnras/stw1450

Publisher’s copyright statement:
This article has been published in Monthly Notices of the Royal Astronomical Society ©: 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society. All rights reserved.

Additional information:

Use policy

The full-text may be used and/or reproduced, and given to third parties in any format or medium, without prior permission or charge, for personal research or study, educational, or not-for-profit purposes provided that:
- a full bibliographic reference is made to the original source
- a link is made to the metadata record in DRO
- the full-text is not changed in any way
The full-text must not be sold in any format or medium without the formal permission of the copyright holders.
Please consult the full DRO policy for further details.
GAMA/H-ATLAS: common star formation rate indicators and their dependence on galaxy physical parameters


Affiliations are listed at the end of the paper

Accepted 2016 June 15. Received 2016 June 15; in original form 2015 November 10

ABSTRACT
We compare common star formation rate (SFR) indicators in the local Universe in the Galaxy and Mass Assembly (GAMA) equatorial fields (∼160 deg²), using ultraviolet (UV) photometry from GALEX, far-infrared and sub-millimetre (sub-mm) photometry from Herschel Astrophysical Terahertz Large Area Survey, and Hα spectroscopy from the GAMA survey. With a high-quality sample of 745 galaxies (median redshift ⟨z⟩ = 0.08), we consider three SFR tracers: UV luminosity corrected for dust attenuation using the UV spectral slope β (SFRUVcorr), Hα line luminosity corrected for dust using the Balmer decrement (BD) (SFRLαcorr), and the combination of UV and infrared (IR) emission (SFRUV+IR). We demonstrate that SFRUVcorr can be reconciled with the other two tracers after applying attenuation corrections by calibrating Infrared excess (IRX; i.e. the IR to UV luminosity ratio) and attenuation in the Hα (derived from BD) against β. However, β, on its own, is very unlikely to be a reliable attenuation indicator. We find that attenuation correction factors depend on parameters such as stellar mass (M*), z and dust temperature (Tdust), but not on Hα equivalent width or Sérsic index.

Key words: methods: statistical – galaxies: photometry – galaxies: statistics – infrared: galaxies – ultraviolet: galaxies.

1 INTRODUCTION
The distribution functions of star formation rates (SFR) at different cosmic epochs (or more commonly, its integrated form, the evolution of the cosmic SFR density) provide fundamental observational tests for theoretical models of galaxy formation and evolution (e.g. Hopkins & Beacom 2006; Madau & Dickinson 2014). Observationally, one can use a variety of indicators at different wavelengths to measure the level of star formation activity in galaxies. The rest-frame ultraviolet (UV) non-ionizing stellar continuum luminosity, where newly formed massive stars emit the bulk of their energy, is often used as a direct SFR indicator, especially at high redshift.

* E-mail: L.Wang@sron.nl

© 2016 The Authors
Published by Oxford University Press on behalf of the Royal Astronomical Society

Hα nebular recombination emission line luminosity, which probes the hydrogen-ionizing photons produced by the most massive and short-lived stars, is another commonly used SFR indicator when spectroscopy is available.

One of the most significant challenges when using UV or Hα line emission as a direct star formation tracer is the effect of dust attenuation, as the process of star formation takes place in dense, cold and often dusty molecular gas clouds. Indeed, some of the most intensely star-forming galaxies are extremely UV-faint, e.g. those selected in the sub-millimetre (sub-mm). To overcome this problem, various empirical or semi-empirical correction methods have been developed to determine the amount of dust attenuation in a galaxy. For example, the power-law spectral slope of the rest-frame UV continuum β (fλ ∝ λβ), or its proxy the far-ultraviolet (FUV)–near-ultraviolet (NUV) colour, has been widely used as...
a practical method for estimating the global attenuation corrections (e.g. Meurer et al. 1995, 1997; Meurer, Heckman & Calzetti 1999; Burgarella, Buat & Iglesias-Páramo 2005; Laird et al. 2005; Reddy et al. 2006; Salim et al. 2007; Treyer et al. 2007; Wijesinghe et al. 2011). However, the spectral slope $\beta$ shows a wide dispersion with varying dust properties, dust/star geometry and redshift (e.g. Granato et al. 2000; Witt & Gordon 2000; Oteo et al. 2014). In addition, $\beta$ is sensitive to the intrinsic UV spectral slope (determined by properties such as the age of the stellar populations, star formation history, metallicity, etc.) and as such is dependent on a number of parameters that are not solely related to dust attenuation (e.g. Kong et al. 2004; Buat et al. 2005).

Another independent method to correct for dust is to use Balmer decrement (BD) ratio measurement (i.e. the observed flux ratio of the $H\alpha$ and $\Hb$ nebular emission lines) to estimate the amount of dust attenuation at $H\alpha$ (e.g. Kennicutt 1992; Brinchmann et al. 2004; Moustakas, Kennicutt & Tremonti 2006; Garn & Best 2010). However, BD measurements are generally only available for bright H II regions within the galaxies, and so can be problematic when applying to the whole galaxy. Also, $\Hb$ is considerably weaker than $H\alpha$. The BD method is also found to be a poor estimator for dust attenuation in dusty starbursts (e.g. Moustakas et al. 2006).

From an energy conservation point of view, one can also derive SFR from dust emission as dust absorbs the UV and optical light from newly formed stars and re-emit predominantly in the far-infrared (FIR) and sub-mm. The main advantage of inferring SFR from the infrared (IR) emission is that it is not affected by dust attenuation. However, some of the IR emission could be caused by heating from the old/evolved stellar populations or active galactic nuclei (AGN) and thus, is not related to recent star formation (e.g. Helou 1986; Popescu et al. 2000; Bell 2003; Natale et al. 2015).

The aim of this paper is to take advantage of the Galaxy and Mass Assembly (GAMA)\(^1\) survey (Driver et al. 2009, 2011; Liske et al. 2015) and associated multiwavelength surveys to carefully examine some of the most commonly used SFR indicators. GAMA provides a large sample of galaxies in the local Universe where photometric information in the UV and IR as well as measurements of key spectral lines such as $H\alpha$ and $\Hb$ are available. The paper is organized as follows. In Section 2, we describe the various surveys and derived data products used in our analysis. In Section 3, we give a brief overview of the three SFR indicators investigated in this paper and our selection criteria in the UV bands, optical emission lines, and IR and sub-mm bands. In Section 4, we study in detail the properties of our galaxy samples selected at different wavelengths and construct a joint UV-$H\alpha$-IR sample. Then focusing on the joint sample, we examine the dust attenuations derived using different methods, and the correlations between various SFR indicators as a function of galaxy physical parameters, such as stellar mass, redshift, UV continuum slope, BD, Infrared excess (IRX; i.e. the total IR to UV luminosity ratio), Sérsic index, $H\alpha$ emission line equivalent width (EW), and dust temperature. Finally, we give conclusions and discussions in Section 5. In an upcoming GAMA paper (Luke et al., in preparation), 12 SFR metrics are examined and calibrated to a mean relation. In this paper, we focus on just three commonly used SFR indicators and inter-compare them using a very high-quality galaxy sample. With random statistical errors minimized, we investigate the influence of different systematic errors on these SFR indicators.

Throughout the paper, we assume a flat cold dark matter cosmological model with $\Omega_M = 0.3$, $\Omega_\Lambda = 0.7$, $H_0 = 70$ km s$^{-1}$ Mpc$^{-1}$.

\(^1\) http://www.gama-survey.org

Flux densities are corrected for Galactic extinction using the $EB - V$ values provided by Schlegel, Finkbeiner & Davis (1998). We use the AB magnitude scale and the Kroupa (Kroupa & Weidner 2003) initial mass function (IMF) unless otherwise stated.

## 2 SAMPLE SELECTION

### 2.1 Spectroscopic and multiwavelength photometric data

GAMA is an optical spectroscopic survey of low-redshift galaxies, mainly conducted at the Anglo-Australian Telescope. GAMA covers three equal-sized fields to an apparent SDSS DR7 (Sloan Digital Sky Survey – Data Release 7) Petrosian r-band magnitude limit of $r = 19.8$ mag: G09, G12 and G15 (centred at a right ascension of $\sim$9, 12, and 14.5 h, respectively), on the celestial equator. We include all GAMA II main survey targets (SURVEY_CLASS $>$=4) with reliable AUTOZ redshifts ($nQ = 3$; Baldry et al. 2014) from the tiling catalogue version 45 (Baldry et al. 2010). We impose a lower redshift limit of 0.01 which is used to remove stars and galaxies for which peculiar motion would overwhelm the Hubble flow, leading to highly uncertain distances based upon recession velocities alone. We also impose an upper redshift limit of 0.5 as there are very few GAMA galaxies above this redshift. After applying these cuts and restricting galaxies to areas with both Galaxy Explorer Mission (GALEX) and Herschel coverage, we are left with a total of 128 170 galaxies which forms our parent sample in this paper.

In addition to the spectroscopic survey, GAMA has also assembled imaging data from a number of independent surveys in order to generate multiwavelength photometric information spanning the wavelength range from 1 nm to 1 mm. Below, we summarize the two main imaging surveys that are relevant to this paper.

#### 2.1.1 GALEX

The GALEX (Martin et al. 2005) conducted a number of major surveys and observer motivated programs, most notably the all-sky imaging survey and the medium imaging survey (MIS). A dedicated programme (the GALEX guest investigator programme GALEX-GAMA) provided further GALEX observations in GAMA fields to MIS depth (1500s). The final collated data provide near-complete NUV and FUV coverage of the primary GAMA II regions. In the three equatorial regions, coverage is at the 90 per cent level in both FUV and NUV bands. The GAMA-GALEX UV catalogue is a combination of archival data and pointed observations on equatorial GAMA fields. The archival data have been used to extend the UV coverage of the GAMA regions as much as possible. In this paper, we use the GAMA GALEX catalogue described in Liske et al. (2015). The full width at half-maximum (FWHM) of the point spread function (PSF) is $\sim 4.2$ and 5.3 arcsec in the FUV and NUV, respectively. As the resolution of the GALEX observations is significantly lower than that of the SDSS data, a variety of methods are employed to address the source blending issue and derive the FUV and NUV fluxes for every GAMA galaxy. For more details, we refer the reader to Liske et al. (2015).

#### 2.1.2 H-ATLAS

The Herschel (Pilbratt et al. 2010) Astrophysical Terahertz Large Area Survey (H-ATLAS) survey (Eales et al. 2010) conducted observations of the three equatorial fields also observed in the
GAMA redshift survey. H-ATLAS images were obtained using Herschel’s fast-scan parallel mode, where the spacecraft is moving at 60 arcsec s⁻¹, with Photometric Array Camera and Spectrometer (PACS; Poglitsch et al. 2010) observing simultaneously at 100 and 160 µm and Spectral and Photometric Imaging Receiver (SPIRE; Griffin et al. 2010) at 250, 350 and 500 µm. The FWHM of the PSF of the telescope is 9, 13, 18, 25 and 35 arcsec at 100, 160, 250, 350 and 500 µm, respectively (Griffin et al. 2010; Ibar et al. 2010). The total area covered by H-ATLAS in the GAMA 09, 12 and 15 fields is approximately 160 deg². Details of the map-making process can be found in Valiante et al. (in preparation).

The PACS and SPIRE photometry are derived from the PACS/SPire maps by measuring the flux in the appropriate optical-IR photometric apertures for a range of aperture sizes, and fitting a simple parametric function (second-order polynomial) to the median flux as a function of aperture area. This contaminating flux level is then subtracted from the target flux to provide a background-corrected flux. The final error in these bands is then half of 1σ quantile range (from 16 per cent to 84 per cent) of the flux as a function of aperture. This implicitly assumes that the dominant error is the subtraction of contaminating flux (see Driver et al. 2016 for more details).

In addition to the forced sub-mm photometry at known optical source positions, in Appendix A, we also consider the blind H-ATLAS source catalogue matched with GAMA galaxies through the likelihood ratio (LR) technique. We show that the use of the LR-matched GAMA/H-ATLAS catalogue does not change our conclusions.

2.2 Properties of GAMA galaxies

The GAMA team has produced a number of catalogues of galaxy physical properties derived from the ammssed multwavelength photometric and spectroscopic data sets. Below, we summarize the properties of the two catalogues which we make use of in this paper.

The stellar masses \( M_{\text{star}} \) in GAMA (Taylor et al. 2011) are estimated based on Spectral energy distribution (SED) fitting of aperture (AUTO) photometry in the rest-frame wavelength range between 3000 and 11 000 Å (approximately \( u \) through \( Y \)) using a grid of synthetic spectra. The spectra are generated by assuming exponentially decaying star formation history, the Bruzual & Charlot (2003) stellar evolution models with the Chabrier (2003) IMF and the Calzetti et al. (2000) dust extinction law. In order to be consistent with SFR calculations in this paper which assumes the Kroupa IMF, we apply a correction to the stellar masses using \( M_{\text{Kroupa}}^{\text{star}} = M_{\text{Chabrier}}^{\text{star}} \times 0.94 \).

A single-Sérsic (1-component only) fit is performed across all passbands in the SDSS \((ugriz)\), the UKIRT Infrared Deep Sky Survey Large Area Survey (Lawrence et al. 2007) \( YJHK \) and VISTA Kilo-degree Infrared Galaxy survey (VIKING) \( YZJHK \) surveys. Galaxy single-Sérsic modelling is achieved using SIGMA v1.0-2 (Structural Investigation of Galaxies via Model Analysis; Kelvin et al. 2012). SIGMA is a wrapper around several contemporary astronomy tools such as Source Extractor (Bertin & Arnouts 1996), PSF Extractor (Bertin & Delorme, private communication) and GALFIT 3 (Peng et al. 2010). GALFIT is the workhorse fitting algorithm used within SIGMA. It uses a simple downhill minimization in order to minimize the global \( \chi^2 \) of the model.

3 DUST-CORRECTED SFR INDICATORS

There are many methods to derive the SFR of a galaxy. Most of them rely on simple linear conversions from luminosities at a given wavelength. In this paper, we consider three SFR indicators widely used in the literature, i.e. UV continuum luminosity corrected for dust attenuation using the UV spectral slope \( \text{SFR}_{\text{UV,corr}} \), H\textalpha{} emission line luminosity corrected for dust attenuation using the BD \( \text{SFR}_{\text{BD,corr}} \), and combination of the escaped and dust reprocessed emission by adding together UV and IR continuum luminosity \( \text{SFR}_{\text{UV+IR,corr}} \). We examine the overall statistical correlations between \( \text{SFR}_{\text{UV+IR,corr}} \) and \( \text{SFR}_{\text{UV,corr}} \) as well as between \( \text{SFR}_{\text{BD,corr}} \) and \( \text{SFR}_{\text{UV+IR,corr}} \). Then, we investigate the ratios between different SFR indicators as a function of various galaxy physical parameters.

3.1 SFR from UV continuum emission

To estimate the dust attenuation correction in the FUV using the observed FUV—NUV colour, we use the empirical relation in Hao et al. (2011)

\[
A_{\text{FUV}} = 3.83 \times [(\text{FUV} - \text{NUV})_{\text{obs}} - 0.022].
\]

This relation is based on calibrating the total IR to FUV luminosity ratio (IRX) and the attenuation in H\textalpha{} line luminosity (derived from BD) against the FUV—NUV colour using a nearby normal star-forming galaxy sample (Moustakas & Kennicutt 2006), which was designed to cover the full range of optical spectral characteristics present in the local galaxy population. For a detailed discussion on the range of applicability of this relation and comparison with other \( A_{\text{FUV}} \)-IRX relations in the literature (e.g. Meurer et al. 1999; Kong et al. 2004; Buat et al. 2005; Burgarella et al. 2005), please refer to Hao et al. (2011). To estimate dust attenuation in the NUV, \( A_{\text{NUV}} \), we use the Calzetti et al. (2000) attenuation law, which yields a relation \( A_{\text{FUV}}/A_{\text{NUV}} = 1.245 \). While Hao et al. (2011) parametrized the dust attenuation using broad-band \textit{GALEX} colour, we will also consider an equivalent parametrization later using the power-law slope of the UV continuum, \( \beta \), which we derive using two methods.

The first method relies on the colour relation from Kong et al. (2004):

\[
\beta = (\log f_{\text{FUV}} - \log f_{\text{NUV}})/\log(\lambda_{\text{FUV}} - \lambda_{\text{NUV}})
\]

with \( f_{\text{FUV}} \) and \( f_{\text{NUV}} \) referring to the monochromatic flux (or flux density) per unit wavelength interval, and effective wavelengths \( \lambda_{\text{FUV}} = 1528 \) Å and \( \lambda_{\text{NUV}} = 2271 \) Å. If using flux densities per unit frequency interval, then \( \beta \) can be derived as

\[
\beta = (\log f_{\text{FUV}} - \log f_{\text{NUV}})/(\log(\lambda_{\text{FUV}} - \lambda_{\text{NUV}})) - 2.
\]

We refer hereafter to this estimate derived using equation (2) or equation (3) as \( \beta_{\text{colour}} \). We can replace \( (\text{FUV} - \text{NUV})_{\text{obs}} \) with \( \beta \) in equation (1) using equation (2) or equation (3):

\[
A_{\text{FUV}} = 1.65 \times \beta + 3.22.
\]

As our galaxy sample spans a considerable range in redshift, it is necessary to apply \( K \)-corrections to the observed FUV—NUV colour, \( (\text{FUV} - \text{NUV})_{\text{obs}} \). \( K \)-corrections as derived from Loveday et al. (2012) using \textit{KCORRECT} v4.2 (Blanton & Roweis 2007) are
applied in the GALEX bands. We also consider another estimate of β by directly fitting the available photometric data with $f_\lambda \propto \lambda^\beta$ (or equivalently $f_\lambda \propto \lambda^{\beta+2}$) in the rest-frame UV range (1200 < $\lambda$ < 3000 Å). This enables us to use all the available data, and the relevant constraints on β at a given redshift. For instance, at $z \gtrsim 0.185$, the SDSS $u$ band (3557 Å) will be used in the power-law fit and above $z = 0.273$, the GALEX FUV band will move out of our fitting window. We refer hereafter to this estimate as $\beta_{\text{fit}}$. In this paper, we prefer $\beta_{\text{fit}}$ over $\beta_{\text{colour}}$ as no K-corrections are needed in the derivation of $\beta_{\text{fit}}$. As a result, $\beta_{\text{fit}}$ should not be affected by the uncertainties in the K-corrections in the UV bands.

Fig. 1 shows the difference in the two β estimates, $\Delta \beta = \beta_{\text{fit}} - \beta_{\text{colour}}$, as a function of redshift. We only show galaxies with signal to noise (S/N) $\geq 5$ in both the NUV and FUV bands and at least two broad-band filters in the rest-frame UV with which to derive $\beta_{\text{fit}}$. Below $z \sim 0.185$, the two estimates are similar to each other but there is a small systematic difference. As the same photometric information is used in deriving $\beta_{\text{fit}}$ and $\beta_{\text{colour}}$ at $z < 0.185$, the small systematic difference is only due to the difference between the UV spectral shape used in KCORRECT and the power-law approximation used in deriving $\beta_{\text{fit}}$. At $z \gtrsim 0.185$, the $u$-band shifts into the rest-frame UV causing further scatter around the systematic difference seen at $z < 0.185$. At $z \gtrsim 0.273$, the FUV band shifts out of the rest-frame UV and so the NUV band is the only common filter used in both types of β estimates, which clearly causes even larger scatter and a large systematic difference between $\beta_{\text{colour}}$ and $\beta_{\text{fit}}$.

To see this more clearly, Fig. 2 shows the normalized histogram in $\Delta \beta$ in three redshift bins for galaxies with at least two broad-band filters in the rest-frame UV. We also investigate the impact of S/N in the FUV and NUV on $\Delta \beta$. The solid histograms correspond to galaxies with S/N $\geq 5$ in FUV and NUV. While the dashed histograms correspond to galaxies with S/N $\geq 10$ in FUV and NUV. In the lowest two redshift bins, i.e. $z = [0.01, 0.185]$ and [0.185, 0.273], there is a small systematic difference between $\beta_{\text{fit}}$ and $\beta_{\text{colour}}$. $\beta_{\text{colour}}$ systematically overestimates the UV continuum slope compared to $\beta_{\text{fit}}$, which means that the inferred dust correction will be systematically higher than using $\beta_{\text{fit}}$. The scatter in $\Delta \beta$ in the medium redshift bin $z = [0.185, 0.273]$ is larger than the lowest redshift bin $z = [0.01, 0.185]$ which is understandable as the $u$ band is used in deriving $\beta_{\text{fit}}$ at $z > 0.185$. The agreement in the two β estimates improves for galaxies selected at higher S/N in the UV bands. In the highest redshift bin, $z = [0.273, 0.5]$, there is a much larger systematic difference as well as a much larger scatter. In Table 1, we list the 16th, 50th, and 84th percentile in $\Delta \beta$ in the three redshift bins, $z = [0.01, 0.185]$, [0.185, 0.273], and [0.273, 0.5]. For galaxies at $z < 0.273$, the median systematic difference between the two β estimates is small (−0.1). In contrast, the median systematic difference in $\Delta \beta$ is very large (>0.5) for galaxies at $z > 0.273$. Therefore, we do not consider those galaxies above $z = 0.273$ any further in this paper.

In the following sections, we define a UV-selected sample by selecting galaxies in the redshift range $z = [0.01, 0.273]$ detected in both FUV and NUV at a S/N of 10 or greater and with at least two broad-band filters in the rest-frame UV with which to derive $\beta_{\text{fit}}$. This UV-selected sample contains a total of 16 920 galaxies, i.e. 19.5 per cent of the parent sample defined in Section 2.1 (after excluding galaxies in the parent sample at $z > 0.273$). The median statistical error (due to measurement error of the flux density) on $\beta$ ($\beta_{\text{fit}}$ or $\beta_{\text{colour}}$) of this sample is 0.2, which is very similar to the nearby star-forming galaxy sample (Moustakas & Kennicutt 2006) studied in Hao et al. (2011). In Section 4.1, we define a joint sample by applying the UV, Hα, and IR selection criteria together. The median statistical error on $\beta$ ($\beta_{\text{fit}}$ or $\beta_{\text{colour}}$) for the joint sample is considerably smaller, 0.1. However, we point out that the statistical error on $\beta_{\text{colour}}$ is likely to be underestimated for our high S/N.
sample, as we have ignored the uncertainties introduced in applying K-corrections in the UV bands.\(^2\)

With GAMA spectroscopic redshifts, we can derive both FUV and NUV luminosities, as \(vL_V\), from the GALEX photometry and correct them for dust attenuation,

\[
L_{UV,\text{corr}} = 10^{0.4vL_V}L_{UV,}, \tag{5}
\]

where UV is either FUV or NUV. Using these dust corrected luminosities, we derive SFR using the calibrations in Hao et al. (2011), Murphy et al. (2011) and summarized Kennicutt & Evans (2012),

\[
\text{SFR}_{\text{FUV,corr}}/M_\odot\text{yr}^{-1} = 1.72 \times 10^{-10}L_{\text{FUV,corr}}/L_\odot, \tag{6}
\]

and

\[
\text{SFR}_{\text{NUV,corr}}/M_\odot\text{yr}^{-1} = 2.60 \times 10^{-10}L_{\text{NUV,corr}}/L_\odot, \tag{7}
\]

which assumes the Kroupa IMF. The statistical errors on the UV-based SFRs are calculated by simply propagating the statistical errors on the FUV and NUV flux and the statistical error on \(\beta_{\text{colour}}\) or \(\beta_{\text{fit}}\).

### 3.2 SFR from H\(_\alpha\) emission line

As outlined in Hopkins et al. (2003) and Gunawardhana et al. (2013), in order to measure SFR for the whole galaxy using the H\(_\alpha\) recombination line luminosity, corrections for the underlying Balmer stellar absorption, dust obscuration as well as the aperture sampled by the fibre are required,

\[
L_{H_\alpha,\text{corr}}(\text{Watts}) = (EW_{H_\alpha} + EW_c) \times 10^{-0.4(M_r - 34.1)}
\times \frac{3 \times 10^{18}}{[6564.61(1 + z)]^{2}} \left( \frac{F_{H_\alpha}/F_{H_\beta}}{2.86} \right)^{2.36}, \tag{8}
\]

where \(EW_{H_\alpha}\) and \(EW_c\) are the positive (i.e. emission) H\(_\alpha\) equivalent width and the constant equivalent width correction, \(M_r\) is the absolute r-band magnitude, \(F_{H_\alpha}/F_{H_\beta}\) is the BD. Briefly, the H\(_\alpha\)-based SFRs used in this investigation are based on \(EW_{H_\alpha}\). A correction for the missing flux due to the size of the fibre aperture is applied to each galaxy using their r-band absolute magnitudes to estimate the continuum luminosity at the wavelength of H\(_\alpha\). As this approach to applying aperture corrections, described in detail in Hopkins et al. (2003), relies on the assumption that the r-band continuum traces the distribution of H\(_\alpha\) emission within a galaxy, the correction can overestimate or underestimate the line luminosity. Based on a large sample of GAMA galaxies with \(z < 0.05\) that has also been observed with the Sydney-AAO Multi-object Integral-field spectrograph instrument (Croom et al. 2012; Bryant et al. 2015), Richards et al. (2016) demonstrated that the SFRs corrected for aperture effects following the method described in equation (8) tend to over estimate SFR by \(\sim 0.1\) dex. This is in agreement of Brough et al. (2013) who found that the SFRs derived based on equation (8) can be overestimated on average by a factor of 1.26, based on observations of 18 GAMA galaxies using the SPIRAL optical integral field unit on the Anglo-Australian Telescope. The aperture effect, of course, reduces with increasing redshift, and the median redshift of our joint UV-H\(_\alpha\)-IR sample defined in Section 4.1 is 0.08

\(^2\)Some studies (e.g. Rosario et al. 2016) do not apply K-corrections in the UV bands as they are expected to be small (<30 per cent out to \(z = 0.15\)). We note that for our high S/N sample, K-corrections actually dominate over pure measurement error. However, as K-corrections in the UV are very uncertain, it is not clear which approach is better.

\(^3\)See Gunawardhana et al. (2013) for a discussion on the impact of the assumption of constant stellar absorption corrections on the H\(_\alpha\) line luminosities.
3.3 SFR from adding UV and IR emission

We can derive the total SFR of a galaxy by adding together the unobscured star formation traced by the observed UV continuum emission and the obscured star formation traced by the IR dust emission. This method is built on an energy balance consideration which argues that all the starlight absorbed at UV and optical wavelengths by interstellar dust grains is re-emitted in the IR and sub-mm (e.g. Popescu et al. 2000; Tuffs et al. 2004; Bianchi 2008; Baes et al. 2010, 2011; De Looze et al. 2012, 2014; Holwerda et al. 2012). However, there are also limitations with this method, e.g. the contribution of the old stellar populations and asymmetric star/dust geometries (e.g. Bell 2003). These effects may average out for a large statistical sample.

The IR luminosity, $L_{\text{IR}}$, is defined as the integrated luminosity from rest-frame 8 to 1000 $\mu$m. To estimate the $L_{\text{IR}}$ of our GAMA galaxies, we have performed SED fitting to WISE 22 $\mu$m and Herschel photometry from 100 to 500 $\mu$m. There are many suites of empirical models and templates that describe the IR SEDs of galaxies. In this paper, we consider five different SED libraries, the Chary & Elbaz (2001; CE) templates, the Dale & Helou (2002; DH) templates, the Lagache, Dale & Puget (2003; LDP) templates, the Smith et al. (2012; DS) templates and the modified blackbody (MBB) templates. The CE library contains 105 template SEDs of different luminosity classes generated to reproduce the observed correlation between mid-IR (MIR) and FIR luminosities (from 7 to 850 $\mu$m) for local galaxies. The DH library contains 64 locally calibrated templates for normal star-forming galaxies, differing in the slope of the power-law distribution of dust mass and dust emissivity as a function of the radiation field intensity. The LDP library contains 92 template spectra of starburst and normal star-forming galaxies, which are constructed and optimised to reproduce statistical quantities like number counts, redshift distributions and the cosmic IR background. The LDP templates for normal star-forming galaxies all have the same shape and are only scaled in luminosity. The DS library is based on a sample of 250 $\mu$m selected galaxies at $z<0.5$ from the H-ATLAS survey. It is worth noting that the CE and DH libraries are constructed from IRAS-selected samples which favour galaxies with a larger warm dust component. In comparison, the H-ATLAS based DS library includes star-forming galaxies that are much colder. For the MBB templates, we follow the parametrization in Hall et al. (2010). We assume that the IR SED is a greybody at low frequencies, and a power law at high frequencies, i.e.

$$\Theta(\nu) = \nu^\beta B_\nu(T_\text{d}), \nu < \nu_0$$

and

$$\Theta(\nu) = \nu^{-\gamma}, \nu \geq \nu_0$$

where $\beta$ is the dust emissivity index and $B_\nu$ is the Planck function with an effective dust temperature $T_\text{d}$. These two functions (equations 10 and 11) are joined at frequencies $\nu_0$ that can be solved from

$$\frac{d \ln [\nu^\beta B_\nu(T_\text{d})]}{d \ln \nu} = -\gamma,$$

with $\gamma = 2$ (Hall et al. 2010). We also fix $\beta = 2$ (Draine & Lee 1984; Mathis & Whiffen 1989).4

4 We have also tried lower $\beta$ values, $\beta = 1.5$ and 1.2 (e.g. Dunne & Eales 2001; Planck Collaboration XVI 2011). We find that the difference in the resulting $L_{\text{IR}}$ is very small. The mean value of log $L_{\text{IR}}(\beta = 2) - \log L_{\text{IR}}(\beta = 1.5)$ and log $L_{\text{IR}}(\beta = 2) - \log L_{\text{IR}}(\beta = 1.2)$ is $-0.03$ and $-0.05$ dex, respectively.

Within each SED library, we select the best-fitting template by finding the one with the lowest $\chi^2$ value, allowing for rescaling of the template. The error on the $L_{\text{IR}}$ derived using the MBB templates is calculated by marginalizing over the effective dust temperature $T_\text{d}$ parameter. The MBB library is the only one out of the five considered here which has a continuous parameter ($T_\text{d}$) characterizing the shape of the SED. The other four libraries (i.e. CE, DH, LDP and DS) have discrete templates which make it difficult to derive marginalized error on the $L_{\text{IR}}$. For this reason and also the fact that the MBB library has a larger IR colour range than the other libraries, we choose the $L_{\text{IR}}$ values derived using the MBB library as our default in the following sections. It is worth pointing out that the lack of polycyclic aromatic hydrocarbons (PAH) emission features in the MBB templates can potentially cause $L_{\text{IR}}$ to be systematically underestimated. Shipley et al. (2013) estimated that the median ratio of PAH luminosity (the sum of emission features at 6.2, 7.7, 8.6, 11.3, 12.7 and 17.0 $\mu$m) to $L_{\text{IR}}$ is 0.09 for a sample of star-forming IR-luminous galaxies. However, the distribution of $L_{\text{PAH}}/L_{\text{IR}}$ is quite wide. In some galaxies, $L_{\text{PAH}}/L_{\text{IR}}$ could be as high as over 20 per cent or as low as 0 per cent.

To match the UV-selected sample in Section 3.1, we construct an IR-selected sample in the redshift range $z = 0.01, 0.273$ by requiring galaxies detected at WISE 22 and Herschel SPIRE 250 $\mu$m at an S/N of 3 or greater. In addition, we require a Herschel PACS detection with S/N $\geq 3$ at 100 or 160 $\mu$m. So, we have a good sampling of the dusty SED in the MIR, close to the peak, and in the Rayleigh–Jeans tail. This IR-selected sample contains a total of 5182 galaxies, i.e. 6.0 per cent of the parent sample defined in Section 2.1 (after excluding galaxies $z > 0.273$). The median statistical error on log $L_{\text{IR}}$ (due to measurement error of the flux densities) for this selected sample is 0.05 dex. For the same sample, the median error on log $L_{\text{IR}}$ is 0.07 dex without the WISE 22 $\mu$m constraint. For the joint UV-Hz-IR sample defined in Section 4.1, the median statistical error on log $L_{\text{IR}}$ is 0.03 dex.

In Fig. 3, we show an example SED fit of a randomly chosen galaxy from the IR-selected sample in the observed frame. The filled dots and 2σ error bars correspond to the measured flux densities and errors at WISE 22 $\mu$m and Herschel wavelengths (100, 160, 250, 350 and 500 $\mu$m). The different coloured lines are the best-fitting SED from different libraries.
SED libraries. Fig. 4 shows the normalized histogram of the difference in log LIR between various libraries for the IR-selected sample. In Table 2, we list the 16th, 50th and 84th percentile in the difference of the total infrared luminosity estimates ∆log LIR using different SED libraries. In the MBB library, we use a default value of β = 2. We also examine ∆log LIR with and without the addition of the WISE 22 μm photometry. The median difference in ∆log LIR is very small in all cases (<0.1 dex).

Table 2. The 16th, 50th and 84th percentile in the difference of the total infrared luminosity estimates ∆log LIR using different SED libraries. In the MBB library, we use a default value of β = 2. We also examine ∆log LIR with and without the addition of the WISE 22 μm photometry. The median difference in ∆log LIR is very small in all cases (<0.1 dex).

<table>
<thead>
<tr>
<th>Library</th>
<th>No 22 μm</th>
<th>With 22 μm</th>
</tr>
</thead>
<tbody>
<tr>
<td>DH – MBB</td>
<td>0.03, 0.08, 0.12</td>
<td>0.07, 0.09, 0.11</td>
</tr>
<tr>
<td>LDP – MBB</td>
<td>-0.05, -0.00, 0.06</td>
<td>-0.06, 0.02, 0.08</td>
</tr>
<tr>
<td>CE – MBB</td>
<td>0.01, 0.05, 0.10</td>
<td>0.04, 0.08, 0.12</td>
</tr>
<tr>
<td>DS – MBB</td>
<td>-0.13, -0.04, 0.03</td>
<td>-0.12, 0.01, 0.09</td>
</tr>
</tbody>
</table>

SEDS. It is clear that the systematic differences in log LIR arising from different SED libraries is relatively small. The median value in ∆log LIR is <0.1 dex with or without adding the WISE 22 μm constraint. As mentioned above, the CE and DH libraries are biased towards galaxies which contain large warm dust content. As a result, the LIR values derived from the CE and DH libraries are systematically higher than the values derived from the MBB library. It is also clear from Table 2 that the lack of PAH features in the MBB templates is not the main cause of the systematic difference between the CE (or the DH) library and the MBB library as the other two libraries (LDP and DS) also have PAH emission features.

Now we can add the obscured star formation traced by IR emission and the unobscured/escaped star formation traced by UV emission to form an estimate of the total SFR of a galaxy. Using the calibrations in Hao et al. (2011) and Kennicutt & Evans (2012), we have

\[
\text{SFR}_{\text{FUV+IR}} / M_\odot \text{yr}^{-1} = 1.72 \times 10^{-10} [L_{\text{FUV}} / L_\odot] + 0.46 \times [L_{\text{IR}} / L_\odot],
\]

(13)

and

\[
\text{SFR}_{\text{NUV+IR}} / M_\odot \text{yr}^{-1} = 2.60 \times 10^{-10} [L_{\text{NUV}} / L_\odot] + 0.27 \times [L_{\text{IR}} / L_\odot],
\]

(14)

Figure 4. The normalized histogram of the difference in the infrared luminosity (in unit of L⊙) for the IR-selected sample estimated from SED fitting using different SED libraries. The dashed line is where there is no systematic difference. The systematic differences in LIR arising from different libraries is relatively small.

Figure 5. Top: the fraction of galaxies selected at different wavelengths (blue filled dots: UV; green asterisks: Hα; red diamonds: IR) compared to the parent sample (defined in Section 2.1) as a function of redshift. Bottom: the fraction of selected galaxies compared to the parent sample as a function of stellar mass over 0.15 < z < 0.273. For the Hα-selected sample, galaxies with 0.15 < z < 0.17 in the parent sample are excluded.

4 CORRELATION BETWEEN DIFFERENT SFR INDICATORS

4.1 Properties of the selected samples

In the top panel in Fig. 5, we compare the fraction of the selected samples at different wavelengths with the parent sample as a function of redshift. The construction of the parent sample is described in Section 2.1. The selection criteria of the UV, Hα and IR samples are summarized in Table 3. The blue dots correspond to galaxies in the UV sample by requiring S/N ≥10 in both FUV and NUV and at least two filters in the rest-frame UV over 0.01 < z < 0.273. The green dots correspond to galaxies in the selected Hα sample by requiring S/N ≥15 in both Hα and Hβ emission lines for all star-forming galaxies over 0.01 < z < 0.273 using the Kauffmann et al. (2003) classification. The red dots correspond to galaxies in the selected IR sample by requiring galaxies 0.01 < z < 0.273 detected at S/N ≥3 at 22 and 250 μm and S/N ≥3 at either 100 or 160 μm. The UV-selected sample contains more objects compared to the IR or the Hα selection. Both the UV and the Hα samples preferentially select lower redshift objects, while the IR-selected sample is better at picking up objects at higher redshifts compared...
to the UV or Hα selection. In the bottom panel in Fig. 5, we show the fraction of the selected samples as a function of stellar mass. It is clear that the UV and Hα-selected samples preferentially select low-mass galaxies (related to the fact that the UV and Hα selection favour galaxies at lower redshifts), while the completeness fraction of the IR-selected sample has a peak in the middle and declines towards both the low-mass and high-mass end. This peak can be explained by limited sensitivity in the IR (so the completeness at low mass, i.e. low SFR, is low) and increasingly higher passive fraction at high masses.

In the following sections, we will use a joint sample by applying the UV, Hα and IR selection criteria together, which contains a total of 745 objects. Fig. 6 shows the full redshift distribution and stellar mass distribution of the UV-Hα-IR joint sample. We can see that the vast majority of our sources are at \( z < 0.15 \). The median redshift of the joint sample is \( z = 0.077 \) and the median stellar mass is \( \log \text{M}_{\text{star}} = 10.13 \text{ M}_\odot \).

### 4.2 Dust attenuation

Before we proceed to examine the correlation between different SFR indicators, we can look at the relation between different estimates of dust attenuation, i.e. the observed FUV–NUV colour or equivalently the UV spectral slope \( \beta \), the total IR to FUV luminosity ratio IRX (i.e. \( \text{IRX} = \log [L_{\text{IR}}/L_{\text{FUV}}] \)), and the attenuation in the Hα line based on the measured BD ratio.

In the top panel of Fig. 7, we show the attenuation in Hα line based on the BD ratio \( A_{\text{Hα}} \) versus the UV spectral slope \( \beta \). The blue line is the bisector fitting to the Moustakas & Kennicutt (2006) sample used in Hao et al. (2011),

\[
A_{\text{Hα}} = 0.8\beta + 1.6, \tag{15}
\]

or equivalently as a function of the observed FUV–NUV colour,

\[
A_{\text{Hα}} = 1.86(F_{\text{FUV}} - N_{\text{NUV}})_{\text{obs}} - 0.04. \tag{16}
\]

Hao et al. (2011) used the correlation between \( \beta \) (or the observed FUV–NUV colour) and \( A_{\text{Hα}} \) to empirically determine the unattenuated power-law slope of the UV continuum or equivalently the intrinsic (dust-free) FUV–NUV colour, \((F_{\text{FUV}} - N_{\text{NUV}})_{\text{int}}\). From equation (16), when \( A_{\text{Hα}} = 0 \), \((F_{\text{FUV}} - N_{\text{NUV}})_{\text{int}} = 0.022\). However, we find that the Hao et al. relation is not a good description of our sample. Instead, the red line is the bisector fitting to our sample in this paper,

\[
A_{\text{Hα}} = 0.98\beta_{\text{fit}} + 1.62, \tag{17}
\]

so, when setting \( A_{\text{Hα}} \) to 0, the dust-free \((F_{\text{FUV}} - N_{\text{NUV}})_{\text{int}} = 0.15 \) for our sample. If using \( \beta_{\text{colour}} \) instead of \( \beta_{\text{fit}} \), then

\[
A_{\text{Hα}} = 0.91\beta_{\text{colour}} + 1.49, \tag{18}
\]

and the dust-free \((F_{\text{FUV}} - N_{\text{NUV}})_{\text{int}} = 0.16 \).

In the bottom panel of Fig. 7, the total IR to FUV luminosity ratio IRX (i.e. \( \text{IRX} = \log [L_{\text{IR}}/L_{\text{FUV}}] \)) is plotted against \( \beta \). The blue line is from Hao et al. (2011),

\[
\text{IRX} = \log \frac{10^{0.4(F_{\text{FUV}} - N_{\text{NUV}})_{\text{int}} - (F_{\text{FUV}} - N_{\text{NUV}})_{\text{obs}}}}{a_{\text{FUV}}} - 1, \tag{19}
\]

where the slope of the UV part of the attenuation curve, \( s_{\text{FUV}} = 3.83 \), and the fraction of the total IR luminosity due to recent star formation, \( a_{\text{FUV}} = 0.46 \). For comparison, the green line is from Boquien et al. (2012) which studied the IRX–\( \beta \) relation on sub-galactic scales in star-forming galaxies selected from the Herschel Reference Survey. The magenta line is the best-fitting IRX–\( \beta \) relation for normal star-forming spirals in the SINGS sample from Muñoz-Mateos et al. (2009). The yellow line is well-known Meurer et al. (1999) relation for starburst galaxies. The cyan line is best-fitting IRX–\( \beta \) relation for nearby starbursts in Kong et al. (2004). The red line is the best fit (of the functional form defined in equation 19) to our sample with \( s_{\text{FUV}} = 3.67 \) and \( a_{\text{FUV}} = 0.46 \). If using \( \beta_{\text{colour}} \) instead of \( \beta_{\text{fit}} \), then we find \( s_{\text{FUV}} = 3.55 \) and \( a_{\text{FUV}} = 0.46 \) for our sample.

Note that in our fitting procedure, we have fixed \( a_{\text{FUV}} \) at 0.46. This is because the correlation between IRX and \( \beta \) is relatively poor, so it

---

**Table 3.** The sample selection criteria in different wavebands. Applying all selection criteria in the UV, Hα and IR leaves us a total sample of 745 objects.

<table>
<thead>
<tr>
<th>Wavelength</th>
<th>Selection criteria</th>
<th>Number of galaxies</th>
</tr>
</thead>
<tbody>
<tr>
<td>UV</td>
<td>( z = [0.01, 0.273], \text{S/N} \geq 10 ) in FUV and NUV, at least two broad-band filters in the rest-frame UV</td>
<td>16,920</td>
</tr>
<tr>
<td>Hα</td>
<td>( z = [0.01, 0.15] ) or ( [0.17, 0.273], \text{S/N} \geq 15 ) in Hα and Hβ, star-forming galaxies</td>
<td>5,717</td>
</tr>
<tr>
<td>IR</td>
<td>( z = [0.01, 0.273], \text{S/N} \geq 3 ) at 250 and 22 ( \mu \text{m}, ) and ( \text{S/N} \geq 3 ) at either 100 or 160 ( \mu \text{m} )</td>
<td>5,182</td>
</tr>
</tbody>
</table>
Figure 7. Top: attenuation in Hα line derived from the Balmer decrement ratio A(Hα) versus UV spectral slope β_Hα for galaxies in our joint UV-Hα-IR sample. The blue dashed line is the bisector fitting used in Hao et al. (2011). The red solid line is the bisector fitting to our sample. The black error bars indicate the median errors of the x- and y-axis. Bottom: IRX (i.e. log 10 L(IR)/L(UV)) versus β_Hα. The blue dashed line is from Hao et al. (2011). The green dotted line is derived for star-forming galaxies in the Herschel Reference Survey from Boquien et al. (2012). The magenta dash--dotted line is derived for the SINGS sample from Muñoz-Mateos et al. (2009). The thick yellow solid line is the relation for starburst galaxies from Meurer et al. (1999). The thick cyan dotted line is the best-fitting relation for starbursts from Kong et al. (2004). The thick red solid line is the best fit to our sample. The black error bars indicate the median errors of the x- and y-axis. The median statistical error on A(Hα), β and IRX is 0.14, 0.12 and 0.04, respectively.

Figure 8. log [L(FUV) + 0.46*L(IR)] versus log [L(NUV) + 0.27*L(IR)] for galaxies in our joint UV-Hα-IR sample. The correlation is very tight. The red line is the expected correlation based on matching the SFR prescriptions (i.e. equations 13 and 14).

4.3 SFR correlations

First of all, we examine how the observed UV luminosities (FUV and NUV) and the Hα line luminosity, both uncorrected for dust attenuation, compare with the linear combination of the observed UV luminosity and the total IR luminosity. Fig. 9 compares the linear combination of the observed UV luminosity and the total IR luminosity with the observed Hα line luminosity.Fig. 10 compares the linear combination of the observed UV luminosity and the total IR luminosity with the observed UV luminosity (without correction for dust attenuation). The solid line in each panel is the predicted relation based on matching SFR prescriptions. As expected, the FUV band suffers more from dust attenuation than the NUV. However, in the luminosity range probed by our sample, we do not see a significant difference in the amount of attenuation with changing UV luminosity. Fig. 10 compares the linear combination of the observed UV luminosity and the total IR luminosity with the observed UV luminosity (without correction for dust attenuation). Again, the solid line in each panel is the predicted relation based on matching SFR prescriptions. The effects of dust obscuration is less severe in the optical emission lines than in the UV. We see a similar amount of attenuation in the observed Hα line luminosity in the top and bottom panel, as expected. Unlike Fig. 9, here we do see some evidence that more actively star-forming galaxies suffer more from dust attenuation.
Now we can examine how different SFR indicators compare after applying corrections for dust attenuation. Starting with the UV SFR indicators, we apply dust correction factors based on the UV spectral slope $\beta$ (or equivalently the observed (FUV − NUV) colour) using equation (4) in Section 3.1. In Fig. 11, we show comparisons between SFR$_{\text{UV, corr}}$ (based on dust-corrected UV luminosity using $\beta$) and SFR$_{\text{UV + IR}}$. There is a good correlation (linear and tight) between SFR$_{\text{UV, corr}}$ and SFR$_{\text{UV + IR}}$, but there is a large overall shift (around 0.3 dex). SFR$_{\text{UV, corr}}$ gives systematically higher values than SFR$_{\text{UV + IR}}$. The offset is likely to be caused by an incorrect $A_{\text{FUV}}−\beta$ relation. We have demonstrated in Section 4.2 that the Hao et al. (2011) $A_{\text{FUV}}−\beta$ relation does not provide a suitable description for our sample. So, we apply the new $A_{\text{FUV}}−\beta$ relation (i.e. equation 22) to derive SFR$_{\text{UV, corr}}$, which are plotted in Fig. 12. It is clear that large offset between SFR$_{\text{UV, corr}}$ and SFR$_{\text{UV + IR}}$ has now gone away. In Table 4, we list the 16th, 50th and 84th percentile in the difference between SFR$_{\text{UV, corr}}$ and SFR$_{\text{UV + IR}}$, where UV means either FUV or NUV. We compare the difference in $\Delta$SFR using $\beta_{\text{fit}}$ or $\beta_{\text{colour}}$. We also compare the difference between applying the Hao et al. (2011) $A_{\text{FUV}}−\beta$ relation and applying the new $A_{\text{FUV}}−\beta$ relations (equations 22 and 23) derived in this paper. The percentiles are similar regardless of using $\beta_{\text{fit}}$ or $\beta_{\text{colour}}$ and regardless of using the FUV or NUV band. There is still a small offset (around 0.1 dex) between SFR$_{\text{NUV, corr}}$ and SFR$_{\text{NUV + IR}}$, using either $\beta_{\text{fit}}$ or $\beta_{\text{colour}}$. This is likely due to the assumption made on the relation between $A_{\text{FUV}}$ and $A_{\text{NUV}}$ (see Section 3.1).

In Fig. 13, we show comparisons between SFR$_{\text{H\alpha, corr}}$ and SFR$_{\text{UV + IR}}$. The dust attenuation corrections applied in SFR$_{\text{H\alpha, corr}}$ are based on the BD measurements. Again, the correlations are similar whether using FUV or NUV. SFR$_{\text{H\alpha, corr}}$ gives systematically higher values than SFR$_{\text{UV + IR}}$. In Table 5, we list the 16th, 50th and 84th percentile in the difference of SFR indicators $\Delta$SFR = SFR$_{\text{UV + IR}}$ − SFR$_{\text{H\alpha, corr}}$. We compare the difference between using the IR luminosity $L_{\text{IR}}$ derived from the MBB template library and using $L_{\text{IR}}$ derived from the Dale & Helou (2002) (DH) template library. It is clear that the small systematic offset between SFR$_{\text{H\alpha, corr}}$ and SFR$_{\text{UV + IR}}$ see in Fig. 13 can be entirely explained.
by systematic error in $L_{\text{IR}}$. The offset could also be due to other systematic errors of a similar scale in the H$\alpha$-based SFR tracer, e.g. optical depth effects in the Balmer corrections, systematic error in Balmer line absorption corrections or aperture effects, etc.

In Appendix C, we further compare the three empirical SFR indicators considered in this paper, i.e. SFR$_{\text{UV, corr}}$, SFR$_{\text{H\alpha, corr}}$ and SFR$_{\text{UV + IR}}$, with SFRs from Grootes et al. (2013) which are derived using the radiative transfer models of Popescu et al. (2011).

4.4 Dependence on galaxy physical parameters

In Fig. 14, we compare the ratios of different SFR indicators as a function of various physical parameters such as stellar mass, redshift, BD, IRX, $\beta_{\text{fit}}$, H$\alpha$ EW, dust temperature and $\text{S}´$ersic index in the SDSS optical bands. In the top panel of each figure, we show the normalized histogram of the x-axis. In the bottom panel of each figure, the blue lines correspond to the 25th, 50th and 75th percentile in the ratio of SFR$_{\text{H\alpha, corr}}$ over SFR$_{\text{UV, corr}}$. The red lines correspond to the 25th, 50th and 75th percentile in the ratio of SFR$_{\text{UV, corr}}$ over SFR$_{\text{UV + IR}}$. In Fig. 15, we show the dust attenuation correction factors applied in SFR$_{\text{H\alpha, corr}}$ and SFR$_{\text{UV, corr}}$ as a function of these physical parameters.

Stellar mass: neither SFR$_{\text{UV, corr}}$/SFR$_{\text{UV + IR}}$ nor SFR$_{\text{H\alpha, corr}}$/SFR$_{\text{UV + IR}}$ has a significant dependence on stellar mass. The dust attenuation corrections applied in SFR$_{\text{UV, corr}}$ and SFR$_{\text{H\alpha, corr}}$ increase with increasing stellar mass indicating more massive galaxies are more obscured. It could also imply that galaxies at higher redshift are more obscured and/or we are preferentially selecting more massive galaxies at higher redshift.

Redshift: neither SFR$_{\text{UV, corr}}$/SFR$_{\text{UV + IR}}$ nor SFR$_{\text{H\alpha, corr}}$/SFR$_{\text{UV + IR}}$ exhibits a significant dependence on redshift. The attenuation corrections applied in SFR$_{\text{UV, corr}}$ and SFR$_{\text{H\alpha, corr}}$ increase with increasing redshift. It could be because galaxies at higher redshift are more obscured and/or we are preferentially selecting more massive galaxies at higher redshift.

BD: the attenuation correction in SFR$_{\text{H\alpha, corr}}$ is uniquely determined by BD. The attenuation correction applied in SFR$_{\text{UV, corr}}$ increases with increasing BD which is expected given the correlation between A(H$\alpha$) and $\beta$ seen in the top panel in Fig. 7. SFR$_{\text{H\alpha, corr}}$/SFR$_{\text{UV + IR}}$ increases with increasing values of BD (i.e. larger attenuation correction) which is most likely due to the fact that the dust-corrected H$\alpha$ line luminosity directly depends on BD.
Table 4. The 16th, 50th and 84th percentile in the difference of SFR indicators $\Delta$SFR = SFR$_{UV, corr}$ − SFR$_{IRMBB, corr}$, where UV means either FUV or NUV. The dust attenuation correction applied to the UV SFR indicator is based on the UV spectral slope $\beta$. We compare the difference in SFR indicators using $\beta_{UV}$ or $\beta_{colour}$. We also compare the difference between applying the Hao et al. (2011) $A_{FUV}$–$\beta$ relation and applying the new $A_{FUV}$–$\beta$ relations (equations 22 and 23) derived in this paper.

<table>
<thead>
<tr>
<th>Hao et al. relation</th>
<th>New relation in this paper</th>
</tr>
</thead>
<tbody>
<tr>
<td>SFR$<em>{FUV+IRMBB}$ − SFR$</em>{UV, corr, \beta_{UV}}$</td>
<td>−0.4, −0.2, 0.0</td>
</tr>
<tr>
<td>SFR$<em>{FUV+IRMBB}$ − SFR$</em>{UV, corr, \beta_{colour}}$</td>
<td>−0.5, −0.3, 0.0</td>
</tr>
<tr>
<td>SFR$<em>{NUV+IRMBB}$ − SFR$</em>{UV, corr, \beta_{UV}}$</td>
<td>−0.4, −0.3, −0.1</td>
</tr>
<tr>
<td>SFR$<em>{NUV+IRMBB}$ − SFR$</em>{UV, corr, \beta_{colour}}$</td>
<td>−0.5, −0.3, −0.1</td>
</tr>
</tbody>
</table>

The SFR$_{FUV, corr}$/SFR$_{FUV+IR}$ ratio decreases with increasing BD, which is caused by the broad correlation between BD and IRX and the dependence of SFR$_{FUV, corr}$/SFR$_{FUV+IR}$ on IRX.

IRX: IRX is a measure of the overall UV photon escape fraction. As expected, the attenuation corrections applied in SFR$_{FUV, corr}$ and SFR$_{H\alpha, corr}$ increase with increasing IRX. This is consistent with Fig. 7. SFR$_{H\alpha, corr}$/SFR$_{FUV+IR}$ shows no appreciable dependence on IRX. SFR$_{FUV, corr}$/SFR$_{FUV+IR}$ decreases significantly with increasing IRX, which is caused by the large scatter in the IRX–$\beta$ correlation shown in the bottom panel in Fig. 7. For objects with high IRX values, it is clear that the dust correction factors $A_{FUV}$ based on the IRX–$\beta$ correlation derived for the whole sample (the red line in the bottom panel in Fig. 7) will underestimate the true level of attenuation.

UV continuum slope $\beta_{UV}$: the attenuation correction applied in SFR$_{FUV, corr}$ is uniquely determined by $\beta$. The attenuation correction applied in SFR$_{H\alpha, corr}$ increases with increasing $\beta$ which is expected given the correlation between $A(H\alpha)$ and $\beta$. Both SFR$_{H\alpha, corr}$/SFR$_{FUV+IR}$ and SFR$_{FUV, corr}$/SFR$_{FUV+IR}$ stays more or less flat with changing $\beta_{UV}$.

H$\alpha$ EW: the H$\alpha$ emission line EW is a measure of specific SFR (Kennicutt, Tamblyn & Congdon 1994). The attenuation correction factors applied in SFR$_{H\alpha, corr}$ and SFR$_{FUV, corr}$ do not vary as a function of the H$\alpha$ EW. SFR$_{FUV, corr}$/SFR$_{FUV+IR}$ stays flat. SFR$_{H\alpha, corr}$/SFR$_{FUV+IR}$ increases with increasing H$\alpha$ EW. One explanation could be that the measurement of the H$\alpha$ EW affects the H$\alpha$ line luminosity.

Dust temperature: the attenuation corrections applied in SFR$_{FUV, corr}$ and SFR$_{H\alpha, corr}$ increase slightly with increasing dust temperature. This is consistent with the fact that galaxies with warmer dust tend to have higher IR luminosities. The SFR$_{H\alpha, corr}$/SFR$_{FUV+IR}$ ratio does not have a significant dependence on dust temperature. However, the SFR$_{FUV, corr}$/SFR$_{FUV+IR}$ ratio decreases with increasing dust temperature which is due to the broad positive correlation between dust temperature and IRX. As explained above, SFR$_{FUV, corr}$/SFR$_{FUV+IR}$ decreases with increasing IRX caused by the large scatter in the IRX–$\beta$ correlation.

Sérsic index: neither the attenuation correction factors or the SFR ratios changes significantly a function of the Sérsic index in the $r$ band. We find similar trends with respect to the Sérsic indexes in the other SDSS optical bands.

5 SUMMARY

In this paper, we compare multiwavelength SFR indicators in the local Universe in the three GAMA equatorial fields. Our analysis uses UV photometry from GALEX, FIR and sub-mm photometry from Herschel H-ATLAS, and H$\alpha$ spectroscopy from the GAMA redshift survey. To minimize random statistical errors, we construct a very high-quality sample of 745 objects (median redshift...
Figure 14. The ratio of SFR indicators as a function of various physical parameters (stellar mass, redshift, Balmer decrement, IRX, $\beta_{\text{IR}}$, H$\alpha$ equivalent width, dust temperature, and Sérsic index in the r band). In the top panel of each figure, we show the normalized histogram of the x-axis. The thin blue lines correspond to the 25th, 50th and 75th percentile in the ratio of SFR$_{\text{H}$\alpha$,\text{corr}}$ over SFR$_{\text{FUV} + \text{IR}}$. The thick red lines correspond to the 25th, 50th and 75th percentile in the ratio of SFR$_{\text{FUV},\text{corr}}$ (using $\beta_{\text{IR}}$ and our new $A_{\text{FUV}}-\beta$ relation equation 22) over SFR$_{\text{FUV} + \text{IR}}$. 

Figure 15. The dust attenuation correction factors as a function of various physical parameters (stellar mass, redshift, Balmer decrement, IRX, $\beta_{H\alpha}$, $H\alpha$ equivalent width, dust temperature, and Sérsic index in the $r$ band). In the top panel of each figure, we show the normalized histogram of the x-axis. The thin blue lines correspond to the 25th, 50th and 75th percentile in the dust correction factors applied in the $H\alpha$-based SFR. The thick red lines correspond to the 25th, 50th and 75th percentile in the dust correction factors applied in the FUV-based SFR.
We find a good linear correlation between SFRUV,corr and SFRUV + IR, but with a ~0.3 dex offset when using the UV spectral slope $\beta$ and the Hao et al. (2011) $A_{\text{FUV}}-\beta$ relation for deriving the dust attenuation correction. This offset is removed when we replace the Hao et al. relation with our new $A_{\text{FUV}}-\beta$ relation based on calibrating IRX and the attenuation in H$\alpha$ against $\beta$. The $A_{\text{FUV}}-\beta$ relation is slightly different depending on whether $\beta_{\text{in}}$ or $\beta_{\text{colour}}$ is used and the choice of IR SED library. The difference between the Hao et al. (2011) $A_{\text{FUV}}-\beta$ relation based on a nearby star-forming sample and the new relation derived in this paper is due to the difference in the galaxy samples. In addition to being at higher redshifts, our galaxy sample corresponds to much lower survey flux limits in the IR and UV and therefore contains many more quiescent star-forming galaxies with redder UV spectral slopes and lower IRX values. We also find a good linear correlation between SFRH$\alpha$,corr and SFRUV + IR. There is a small median offset of around 0.1 dex. But we demonstrate that this offset can be entirely explained by systematic effects in deriving the IR luminosity, $L_{\text{IR}}$, and/or other systematic errors in the H$\alpha$-based SFR tracer. Moreover, the correlation between SFRH$\alpha$,corr and SFRUV + IR has a similar scatter (0.2 dex) as the correlation between SFRUV,corr and SFRUV + IR.

The ratios between different SFR indicators and the dust attenuation correction factors applied in the UV (using $\beta$) and H$\alpha$ (using the BD) are examined as a function of various galaxy physical parameters. The attenuation factor applied in SFRH$\alpha$,corr, which is uniquely determined by BD increases with increasing values of IRX and $\beta$. Similarly, the attenuation factor applied in SFRUV, $\beta$ which is uniquely determined by $\beta$ increases with increasing values of BD and IRX. These trends are consistent with the broad correlations between BD, $\beta$, and IRX seen in Fig. 7. We also find that attenuation correction factors depend on stellar mass, redshift and dust temperature, but not on the H$\alpha$ EW or Sérsic index in the SDSS optical bands.

After applying corrections for dust attenuation, we find that the SFRUV,corr/SFRUV + IR ratio does not depend significantly on stellar mass, redshift, UV spectral slope $\beta$, H$\alpha$ EW, or structural parameters such as Sérsic index. However, the SFRUV,corr/SFRUV + IR ratio does systematically decrease with increasing values of IRX, BD, and dust temperature. The dependence on IRX is caused by the large scatter in the IRX versus $\beta$ relation. For objects with high IRX values, the dust attenuation correction factor $A_{\text{FUV}}$ based on the IRX–$\beta$ correlation derived for the whole sample will underestimate the true level of attenuation. Also, there is a positive correlation between IRX and BD and between IRX and dust temperature which explains the systematic trend in the SFRUV,corr/SFRUV + IR ratio as a function of BD and dust temperature. In contrast, the SFRH$\alpha$,corr/SFRUV + IR ratio does not show any systematic trend as a function of various physical parameters except BD and H$\alpha$ EW, which is most likely caused by the fact that both Balmer decrement and H$\alpha$ EW directly determine the dust-corrected H$\alpha$ line luminosity.

**ACKNOWLEDGEMENTS**

We thank the anonymous referee whose comments improved the paper significantly. LW acknowledges support from an ERC StG grant (DEGAS-259586). PN acknowledges the support of the Royal Society through the award of a University Research Fellowship, the European Research Council, through receipt of a Starting Grant (DEGAS-259586) and support of the Science and Technology Facilities Council (ST/L00075X/1). SB acknowledges the funding support from the Australian Research Council through a Future Fellowship (FT140101166). NB acknowledges the support of the EC FP7 SPACE project ASTRODEEP (Ref. No. 312725). RJJ, LD, SJM and IO acknowledge support from ERC in the form of the Advanced Investigator Programme, 321302, COSMICISM. E. Ibar acknowledges funding from CONICYT/FONDECYT postdoctoral project N°3130504. MALL acknowledges support from UNAM through the PAPIIT project IA101315. SD acknowledges support from a STFC Ernest Rutherford fellowship.

GAMA is a joint European-Australasian project based around a spectroscopic campaign using the Anglo-Australian Telescope. The GAMA input catalogue is based on data taken from the SDSS and the UKIRT Infrared Deep Sky Survey. Complementary imaging of the GAMA regions is being obtained by a number of independent survey programs including GALEX MIS, VST KiDS, VISTA VIKING, WISE, Herschel-ATLAS, GMRT and ASKAP providing UV to radio coverage. GAMA is funded by the STFC (UK), the ARC (Australia), the AAO, and the participating institutions. The GAMA website is http://www.gama-survey.org/.

The H-ATLAS is a project with Herschel, which is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA. The H-ATLAS web site is http://www.h-atlas.org/.

We acknowledge the use of data products from the NASA-operated GALEX space mission. We also acknowledge the following institutions and agencies for their financial contributions towards the reactivation and operations of the GALEX satellite. This has allowed us to complete NUV observations of the G23 and Herschel ATLAS SGP regions: The Australian-Astronomical Observatory (AOO), the Australian Research Council (ARC), the National Centre for Radio Astronomy Research (ICRAR), the University of Western Australia, the University of Sydney, the University of Canterbury, Max Planck Institute fuer Kernphysik (MPIK), the University of Queensland, the University of Edinburgh, Durham University, the European Southern Observatory (ESO), the University of Central Lancashire, Liverpool John Moore University, National Aeronautics and Space Administration (NASA), Université Paris Sud, University of California Irvine, Instituto Nazionale Di Astrofisica (INAF), and the University of Hertfordshire.

**REFERENCES**


MNRA461, 1898–1916 (2016)
APPENDIX A: H-ATLAS BLIND CATALOGUE

The large PSF of Herschel imaging and the relatively poor correlation between sub-mm and optical brightness means that matching sub-mm sources to optical sources can be ambiguous. The two techniques considered represent alternative solutions to the problem: the technique used in the main body of the paper is forced matching between different SFR indicators using the H-ATLAS blind source catalogue. Therefore, it is important to test whether the correlations we see between SFR_{UV}+IR and SFR_{UV,corr} and between SFR_{UV}+IR+ and SFR_{Hα,corr} depend on how Herschel fluxes are derived for GAMA galaxies. Here we re-derive the correlations between different SFR indicators using the H-ATLAS blind source catalogue.

The H-ATLAS catalogue contains all SPIRE sources which are identified on the LR method (e.g. Sutherland & Saunders 1992; Wang & Rowan-Robinson 2009) and those with ambiguous optical counterparts. Therefore, it is important to test whether the correlations we see between SFR_{UV}+IR and SFR_{UV,corr} and between SFR_{UV}+IR+ and SFR_{Hα,corr} depend on how Herschel fluxes are derived for GAMA galaxies. Here we re-derive the correlations between different SFR indicators using the H-ATLAS blind source catalogue.

The large PSF of Herschel imaging and the relatively poor correlation between sub-mm and optical brightness means that matching sub-mm sources to optical sources can be ambiguous. The two techniques considered represent alternative solutions to the problem: the technique used in the main body of the paper is forced matching between different SFR indicators using the H-ATLAS blind source catalogue. Therefore, it is important to test whether the correlations we see between SFR_{UV}+IR and SFR_{UV,corr} and between SFR_{UV}+IR+ and SFR_{Hα,corr} depend on how Herschel fluxes are derived for GAMA galaxies. Here we re-derive the correlations between different SFR indicators using the H-ATLAS blind source catalogue.

The H-ATLAS catalogue contains all SPIRE sources which are identified on the LR method (e.g. Sutherland & Saunders 1992; Wang & Rowan-Robinson 2009) and those with ambiguous optical counterparts. Therefore, it is important to test whether the correlations we see between SFR_{UV}+IR and SFR_{UV,corr} and between SFR_{UV}+IR+ and SFR_{Hα,corr} depend on how Herschel fluxes are derived for GAMA galaxies. Here we re-derive the correlations between different SFR indicators using the H-ATLAS blind source catalogue.

The large PSF of Herschel imaging and the relatively poor correlation between sub-mm and optical brightness means that matching sub-mm sources to optical sources can be ambiguous. The two techniques considered represent alternative solutions to the problem: the technique used in the main body of the paper is forced matching between different SFR indicators using the H-ATLAS blind source catalogue. Therefore, it is important to test whether the correlations we see between SFR_{UV}+IR and SFR_{UV,corr} and between SFR_{UV}+IR+ and SFR_{Hα,corr} depend on how Herschel fluxes are derived for GAMA galaxies. Here we re-derive the correlations between different SFR indicators using the H-ATLAS blind source catalogue.

The large PSF of Herschel imaging and the relatively poor correlation between sub-mm and optical brightness means that matching sub-mm sources to optical sources can be ambiguous. The two techniques considered represent alternative solutions to the problem: the technique used in the main body of the paper is forced matching between different SFR indicators using the H-ATLAS blind source catalogue. Therefore, it is important to test whether the correlations we see between SFR_{UV}+IR and SFR_{UV,corr} and between SFR_{UV}+IR+ and SFR_{Hα,corr} depend on how Herschel fluxes are derived for GAMA galaxies. Here we re-derive the correlations between different SFR indicators using the H-ATLAS blind source catalogue.

The large PSF of Herschel imaging and the relatively poor correlation between sub-mm and optical brightness means that matching sub-mm sources to optical sources can be ambiguous. The two techniques considered represent alternative solutions to the problem: the technique used in the main body of the paper is forced matching between different SFR indicators using the H-ATLAS blind source catalogue. Therefore, it is important to test whether the correlations we see between SFR_{UV}+IR and SFR_{UV,corr} and between SFR_{UV}+IR+ and SFR_{Hα,corr} depend on how Herschel fluxes are derived for GAMA galaxies. Here we re-derive the correlations between different SFR indicators using the H-ATLAS blind source catalogue.

The large PSF of Herschel imaging and the relatively poor correlation between sub-mm and optical brightness means that matching sub-mm sources to optical sources can be ambiguous. The two techniques considered represent alternative solutions to the problem: the technique used in the main body of the paper is forced matching between different SFR indicators using the H-ATLAS blind source catalogue. Therefore, it is important to test whether the correlations we see between SFR_{UV}+IR and SFR_{UV,corr} and between SFR_{UV}+IR+ and SFR_{Hα,corr} depend on how Herschel fluxes are derived for GAMA galaxies. Here we re-derive the correlations between different SFR indicators using the H-ATLAS blind source catalogue.

The large PSF of Herschel imaging and the relatively poor correlation between sub-mm and optical brightness means that matching sub-mm sources to optical sources can be ambiguous. The two techniques considered represent alternative solutions to the problem: the technique used in the main body of the paper is forced matching between different SFR indicators using the H-ATLAS blind source catalogue. Therefore, it is important to test whether the correlations we see between SFR_{UV}+IR and SFR_{UV,corr} and between SFR_{UV}+IR+ and SFR_{Hα,corr} depend on how Herschel fluxes are derived for GAMA galaxies. Here we re-derive the correlations between different SFR indicators using the H-ATLAS blind source catalogue.

The large PSF of Herschel imaging and the relatively poor correlation between sub-mm and optical brightness means that matching sub-mm sources to optical sources can be ambiguous. The two techniques considered represent alternative solutions to the problem: the technique used in the main body of the paper is forced matching between different SFR indicators using the H-ATLAS blind source catalogue. Therefore, it is important to test whether the correlations we see between SFR_{UV}+IR and SFR_{UV,corr} and between SFR_{UV}+IR+ and SFR_{Hα,corr} depend on how Herschel fluxes are derived for GAMA galaxies. Here we re-derive the correlations between different SFR indicators using the H-ATLAS blind source catalogue.

The large PSF of Herschel imaging and the relatively poor correlation between sub-mm and optical brightness means that matching sub-mm sources to optical sources can be ambiguous. The two techniques considered represent alternative solutions to the problem: the technique used in the main body of the paper is forced matching between different SFR indicators using the H-ATLAS blind source catalogue. Therefore, it is important to test whether the correlations we see between SFR_{UV}+IR and SFR_{UV,corr} and between SFR_{UV}+IR+ and SFR_{Hα,corr} depend on how Herschel fluxes are derived for GAMA galaxies. Here we re-derive the correlations between different SFR indicators using the H-ATLAS blind source catalogue.

The large PSF of Herschel imaging and the relatively poor correlation between sub-mm and optical brightness means that matching sub-mm sources to optical sources can be ambiguous. The two techniques considered represent alternative solutions to the problem: the technique used in the main body of the paper is forced matching between different SFR indicators using the H-ATLAS blind source catalogue. Therefore, it is important to test whether the correlations we see between SFR_{UV}+IR and SFR_{UV,corr} and between SFR_{UV}+IR+ and SFR_{Hα,corr} depend on how Herschel fluxes are derived for GAMA galaxies. Here we re-derive the correlations between different SFR indicators using the H-ATLAS blind source catalogue.
Figure A1. Comparison of $\log_{10} SFR_{\text{FUV}}(\text{corr}$ (using $\beta_{\text{fit}}$) with $\log_{10} SFR_{\text{FUV}} + \text{IR}$ (left-hand panel) and $\log_{10} SFR_{\text{H}\alpha,\text{corr}}$ with $\log_{10} SFR_{\text{FUV}} + \text{IR}$ (right-hand panel) (colour-coded by galaxy counts). The H-ATLAS catalogue is matched to the GAMA galaxies using the likelihood ratio (LR) technique. The black line is the one-to-one relation.

For more details, we refer the reader to Smith et al. (2011) and Bourne et al. (in preparation).

In Fig. A1, we show the correlations between different SFR indicators, using the H-ATLAS blind source catalogue matched with GAMA galaxies through the LR method. The resulting correlations are very similar to Figs 12 and 13 in Section 4.3. Note that we have applied our new $A_{\text{FUV}} - \beta_{\text{int}}$ relation (i.e. Eq 22) and not the Hao et al. (2011) relation.

APPENDIX B: COMPARISON WITH THE HAO ET AL. (2011) SAMPLE AND RELATION

To understand the differences seen in the IRX versus $\beta$ relation between Hao et al. (2011) and this paper, we compare in detail the galaxy samples used in both studies.

Fig. B1 shows IRX versus $\beta_{\text{colour}}$. To better compare Hao et al. (2011) and this paper, $\beta_{\text{colour}}$ is used in this figure. The blue stars correspond to the MK06 galaxy sample studied in Hao et al. (2011) and the blue dashed line is the best fit from Hao et al. (2011). The black symbols correspond to our joint UV-H$\alpha$-IR sample of 745 objects using IR luminosities measured from SED fitting to the MBB library. The black dotted line is our best-fitting IRX-$\beta$ relation to the black symbols. To better match the IR luminosity measurement from Hao et al. (2011), we also plot our measurement using IR luminosities estimated from SED fitting to the Dale & Helou (2002) library which are shown as the red symbols. The red solid line is the best fit to the red symbols.

Fig. B2. Top: infrared luminosity $L_{\text{IR}}$ (in unit of $L_{\odot}$) versus comoving distance. The blue stars correspond to the MK06 galaxy sample. The red symbols correspond to our joint UV-H$\alpha$-IR sample with $L_{\text{IR}}$ estimated from SED fitting to the Dale & Helou (2002) library. Bottom: similar to the top panel but with the y-axis replaced by the observed FUV luminosity $L_{\text{FUV}}$ uncorrected for dust attenuation.

Fig. B1 shows IRX versus $\beta$ using the Moustakas & Kennicutt (2006; hereafter MK06) galaxy sample studied in Hao et al. (2011) and our joint UV-H$\alpha$-IR sample of 745 objects. To better compare the Hao et al. (2011) study and this paper, $\beta_{\text{colour}}$ is used in this figure instead of $\beta_{\text{int}}$. The blue stars correspond to the MK06 sample and...
SFR indicators and galaxy physical parameters

the blue line is the best fit of the functional form defined in equation (19) from Hao et al. (2011) with $s_{\text{FUV}} = 3.83$ and $a_{\text{FUV}} = 0.46$. The black symbols correspond to our joint UV-Hα-IR sample of 745 objects using IR luminosities measured from SED fitting to the MBB library. The black line is our best-fitting IRX–β relation to the black symbols with $s_{\text{FUV}} = 3.55$ and $a_{\text{FUV}} = 0.46$. To better match the IR luminosity measurement from Hao et al. (2011), we also plot our measurement using IR luminosities estimated from SED fitting to the Dale & Helou (2002) library which are shown as the red symbols. The red line is the best fit to the red symbols with $s_{\text{FUV}} = 3.66$ and $a_{\text{FUV}} = 0.46$. As discussed in Section 3.3, $L_{\text{IR}}$ estimated using the DH library is systematically higher than $L_{\text{IR}}$ estimated from the MBB library. However, the median difference is less than 0.1 dex between the two libraries. Comparing the blue stars with the red symbols, it is clear that our joint UV-Hα-IR sample has a lot more quiescent star-forming galaxies with redder UV spectra and lower IRX values.

In Fig. B2, we compare bolometric $L_{\text{IR}}$ versus coming distance and the observed FUV luminosity $L_{\text{FUV}}$ (without correction for dust) versus coming distance between the MK06 sample and our joint UV-Hα-IR sample. The galaxies in our sample are at much higher redshifts than the MK06 sample. In addition, our galaxies correspond to a much lower survey flux limit than the MK06 sample. This is mostly likely due to the fact that the MK06 sample are restricted to galaxies which are detected by IRAS at 25, 60 and 100 μm. As such, the MK06 sample is biased towards warmer dust temperature and more IR-luminous galaxies than our galaxy sample.

APPENDIX C: COMPARISON WITH SFR DERIVED FROM RADIATIVE TRANSFER MODELLING

Arguably an accurate determination of SFR requires radiative transfer modelling of the panchromatic SEDs of galaxies, which could then be used to calibrate SFRs derived from other SFR indicators. In Fig. C1, we compare our SFR indicators ($\text{SFR}_{\text{UV}+\text{IRMBB}}$, $\text{SFR}_{\text{UV, corr}}$ and $\text{SFR}_{\text{Hα, corr}}$) with the NUV-based SFRs (Grootes et al. 2013) derived using the radiative transfer models of Popescu et al. (2011) for a sample of local GAMA spiral galaxies at $z < 0.13$, $\text{SFR}_{\text{RT}}$. The black line indicates the one-to-one relation. In Table B1, we list the 16th, 50th and 84th percentile in the difference between the SFR indicators ($\text{SFR}_{\text{UV}+\text{IRMBB}}$, $\text{SFR}_{\text{UV, corr}}$ and $\text{SFR}_{\text{Hα, corr}}$) studied in this paper and the radiation transfer-corrected NUV-based SFR$_{\text{RT}}$. There is a small median difference in all cases except in $\text{SFR}_{\text{UV, corr}} - \text{SFR}_{\text{RT}}$, which could be because the radiation transfer corrected SFR is also derived from the observed NUV luminosity. Finally, the correlation between $\text{SFR}_{\text{Hα, corr}}$ and $\text{SFR}_{\text{RT}}$ has considerably larger scatter compared to the correlations seen between $\text{SFR}_{\text{UV}+\text{IRMBB}}$ and $\text{SFR}_{\text{RT}}$ and between $\text{SFR}_{\text{UV, corr}}$ and $\text{SFR}_{\text{RT}}$.

### Table C1

<table>
<thead>
<tr>
<th>SFR Indicator</th>
<th>16th Percentile</th>
<th>Median</th>
<th>84th Percentile</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\text{SFR}<em>{\text{UV}+\text{IRMBB}} - \text{SFR}</em>{\text{RT}}$</td>
<td>$-0.3$, $-0.2$, $0.0$</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\text{SFR}<em>{\text{UV, corr}} - \text{SFR}</em>{\text{RT}}$</td>
<td>$-0.2$, $-0.1$, $0.0$</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\text{SFR}<em>{\text{Hα, corr}} - \text{SFR}</em>{\text{RT}}$</td>
<td>$-0.3$, $-0.1$, $0.0$</td>
<td>$-0.1$, $0.0$, $0.2$</td>
<td></td>
</tr>
</tbody>
</table>

1. SRON Netherlands Institute for Space Research, Landleven 12, NL-9747 AD Groningen, the Netherlands
2. Kapteyn Astronomical Institute, University of Groningen, Postbus 800, NL-9700 AV Groningen, the Netherlands
3. Department of Physics, Institute for Computational Cosmology, Durham University, Durham DH1 3LE, UK
4. Department of Astronomy, University of Maryland, College Park, MD 20742-2421, USA