Air/Liquid Interfacial Nanoassembly of Molecular Building Blocks into Preferentially Oriented Porous Organic Nanosheet Crystals via Hydrogen Bonding

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Supporting Information

ABSTRACT: Nanosheets with highly regulated nanopores are ultimately thin functional materials for diverse applications including molecular separation and detection, catalysis, and energy conversion and storage. However, their availability has hitherto been restricted to layered parent materials, covalently bonded sheets, which are layered via relatively weak electrostatic interactions. Here, we report a rational bottom-up methodology that enables nanosheet creation beyond the layered systems. We employ the air/liquid interface to assemble a triphenylbenzene derivative into perfectly oriented highly crystalline noncovalent-bonded organic nanosheets under ambient conditions. Each molecular building unit connects laterally by hydrogen bonding, endowing the nanosheets with size- and position-regulated permanent nanoporosity, as established by in situ synchrotron X-ray surface crystallography and gas sorption measurements. Notably, the nanosheets are constructed specifically by interfacial synthesis, which suppresses the intrinsic complex interpenetrated structure of the bulk crystal. Moreover, they possess exceptional long-term and thermal stability and are easily transferrable to numerous substrates without loss of structural integrity. Our work shows the power of interfacial synthesis using a suitably chosen molecular component to create two-dimensional (2D) nanoassemblies not accessible by conventional bulk crystal exfoliation techniques.

KEYWORDS: nanosheet, air/liquid interface, porous material, membrane, self-assembly, two-dimensional material, oriented film
ultrathin components in nanotechnological devices.\textsuperscript{1,2} Incorporation of highly regulated nanopores into the 2D structure leads to additional enormous capacity for diverse applications.\textsuperscript{3−5} Nanosheets so far are greatly restricted to layered or lamellar systems in which individual sheets are formed by strong interactions between atoms/molecules such as covalent or coordination bonds, and the sheets stack via relatively weak electrostatic interactions.\textsuperscript{6} By utilizing this hierarchy of interactions, layered bulk materials can be delaminated by chemical/physical treatments, leading to the isolation of nanosheets.\textsuperscript{7−9} Nanosheets with well-defined nanopores have been created with layered zeolites\textsuperscript{10,11} and metal− and covalent−organic frameworks\textsuperscript{12,13} by the delamination process. Such top-down exfoliation not only limits the variety of materials but also restricts their use because it necessitates multiproducting, which includes bulk synthesis, exfoliation, stabilization in solution, and deposition on a substrate, while it still remains difficult to place them uniformly on aimed positions.\textsuperscript{7} Although bottom-up approaches utilizing liquid interfaces for creating porous nanosheets have been reported, applicability is still limited to covalently bonded layered systems, including coordination networks.\textsuperscript{14−18} On the other hand, noncovalent bonds/interactions such as van der Waals interactions, dipolar coupling, and hydrogen bonding also play important roles for creating a variety of supramolecular assemblies, including 2D networks.\textsuperscript{19,20}

Developing a rational methodology for creating nanosheets beyond such layered systems is a critical issue to be addressed that will further open the way for their use in a range of important applications. Indeed, there are various molecular-based functional solids, including biomaterials, which assemble by utilizing only electrostatic interactions such as hydrogen bonding and van der Waals forces.\textsuperscript{21,22} One of the important features of these
materials is that they can assemble under mild conditions without recourse to any high-temperature treatment. On the other hand, in forming stable structures, they tend to pack densely to maximize the attractive interactions and minimize empty space. Therefore, such materials tend to adopt nonporous structures.23 As a result, even for bulk systems, creating molecular-based porous frameworks is an ongoing target24−26 and obtaining such porous systems as nanosheets constitutes an even greater challenge.

Here, we present a quick operationally facile direct bottom-up procedure of assembling highly crystalline organic nanosheets without any intermolecular covalent bonds utilizing air/liquid interfaces under ambient conditions (Figure 1a top). The nanosheets are endowed with both size- and position-regulated nanopores and are completely oriented—the pore channels align vertically to the liquid surface. They are easily transferable to various substrates, including holey grids, while retaining their highly crystalline structure (Figure 1b). Notably, the crystal structure of the nanosheets differs from that of their equivalent bulk crystals, which adopt a multiply interpenetrated structure (Figure 1a (bottom),c). The highly crystalline porous 2D network is specifically realized at the air/liquid interface and is not accessible by conventional film fabrication techniques such as dip casting (Figure 1d), drop casting (Figure 1e), and vacuum deposition (Figure 1f, see the Methods). Our work shows the power of interfacial synthesis using molecular building blocks for creating 2D nanomaterials.

RESULTS AND DISCUSSION

Nanosheet Preparation. In order to assemble the porous organic nanosheets, we selected as a building block 1,3,5-tris-(4-carboxyphenyl)-benzene (BTB), which possesses peripheral carboxylic acid linkers (Figure 1). Nanosheet (LINAS-1: liquid-interfacially assembled nanosheet no. 1) growth involved spreading...
of BTB solution in toluene/methanol onto a water subphase in a Langmuir trough at room temperature (Figure 1b). Nanosheet formation at the air/liquid surface was followed by monitoring the surface pressure–mean molecular area ($\Pi$–$A$) isotherms (Figure 2a) and by Brewster angle microscopy (BAM) as the surface was compressed (Figure 2a (inset) and Figure S1).

**Structural Characterization of LINAS-1 by in Situ GIXRD Measurements at the Air/Liquid Interface.** Detailed insight into the LINAS-1 formation and its molecular order was obtained by *in situ* grazing-incidence synchrotron X-ray diffraction (GIXRD) measurements directly at the air/liquid interface. The observation of sharp clearly resolved Bragg peaks in the in-plane (sensitive to the lattice dimensions parallel to the interface) GIXRD profile of LINAS-1 at $\Pi = 0$ mN/m (Figure 2b) provides the signature that assembly of a large-scale highly crystalline structure occurs immediately upon spreading BTB onto the water surface. All peaks to a scattering vector, $Q_{xy} = 2.1 \text{ Å}^{-1}$ index as ($h \kappa 0$) on a metrically hexagonal unit cell with in-plane lattice parameters, $a = b = 31.17(1) \text{ Å}$. Out-of-plane GIXRD patterns (Figure 2c) reveal that LINAS-1 is also highly ordered in the direction vertical to the interface. Increasing $\Pi$ leads to an increase in peak intensity without a change in peak position for both in-plane and out-of-plane GIXRD measurements implying that surface compression does not affect the highly crystalline molecular organization but only the water surface coverage (Figure 2b,c).

Investigation of the crystalline LINAS-1 structure was initiated by considering as starting model the bulk BTB crystal structure, which we determined from single-crystal XRD data (Figures S2–S6 and Table S1). BTB crystallizes with a monoclinic structure comprising interpenetrated 2D honeycomb networks. Hexagonal networks of size, $a = b \sim 31.2$ Å are constructed by hydrogen-bonded carboxylic acid dimers (Figure 1a and Figure S5) and form columns via BTB $\pi$–$\pi$ stacking in an eclipsed-slipped arrangement (slip angle, $\delta = 1.6^\circ$–$5.2^\circ$) at interlayer distances of 3.6–3.8 Å (Figure S6). Three- or four-BTB slabs align along different directions resulting in a complex interpenetrated framework structure (Figure 1a) comparable with those of other porous materials including BTB polymorphs.

The LINAS-1 structural model (Figure 2d) incorporates the same 2D honeycomb motif of hydrogen-bond–linked BTB molecules. This is supported by the comparable in-plane lattice sizes of LINAS-1 and bulk BTB together with evidence from IR spectroscopy (Figure S7). However, the LINAS-1 in-plane XRD
profile is different from that of bulk BTB (Figures S8 and S9), and none of the Bragg reflections coincide with those measured in out-of-plane geometry (Figure 2b). This implies that LINAS-1 is free from the interpenetrated motif of bulk BTB and the honeycomb network plane is perfectly oriented parallel to the water surface. With the aid of the liquid interface as a crafting medium, exquisite control of the molecular orientation is achieved at the nanoscale and the formation of the complex interpenetrated structure is suppressed. The derived interlayer distance of 3.6 Å—consistent with that obtained from single crystal XRD also naturally points to the adoption of a π−π-stacked structure vertical to the water surface (Figure 2e). In-plane XRD profile simulations using eclipsed AA and staggered AB layer-stacking modes clearly favor the former (Figures S10−S12). Allowing for a slip angle (δ = 5°) in the eclipsed stacking motif of the layers leads to best agreement with experiment (Figure 2f).

Integrity and Morphological Characterization of LINAS-1 after Transfer onto Solid Substrates. LINAS-1 is remarkably versatile and exceptionally robust. First, it can be easily transferred to a variety of solid substrate surfaces including holey supports (Figure 1a) and grow in a layer-by-layer manner by a horizontal dipping method (Figure 3a) without loss of structural integrity. This is immediately apparent by comparison of in-plane and out-of-plane synchrotron XRD patterns for LINAS-1 deposited on a Si(100) substrate and LINAS-1 at the air/liquid interface (II = 5 mN/m, Figure 3b,c). There are no differences in position and relative intensity of the Bragg reflections (a change in peak width results from differences in instrumental resolution, see the Methods). The results confirm that the preferentially oriented highly crystalline LINAS-1 structure is retained intact after transfer on the Si substrate.

The sheet thickness and surface morphology of LINAS-1 transferred on Si were evaluated by atomic force microscopy (AFM) (Figure 3d). The height (thickness, Figure 3e) distribution reveals monodispersity with lateral dimension of the nanosheets (maximum Feret diameter, Figure S13) of 300 ± 10 nm. In particular, LINAS-1 has excellent height uniformity with average height of 9.3(2) nm, implying an aspect ratio greater than 30, in good agreement with the crystalline domain size, ~9 nm estimated from the full-width-at-half-maximum (fwhm) of the out-of-plane peak (Figure 3c) using Scherrer’s equation. Given the interlayer distance of 3.6 Å, this implies that ~25 BTB molecules stack along the out-of-plane direction uniformly throughout the nanosheet. The Feret diameter also compares well with the lateral crystalline domain size (~150 nm) estimated from the in-plane GIXRD results. This implies that each nanosheet observed in the AFM measurements corresponds to a single crystalline domain. A red-shifted peak observed in the photoluminescence (PL) spectrum of LINAS-1 supports the long-range stacking order compared to the evaporated amorphous film and the bulk crystal in which BTB stacking is limited to four molecules (Figure S14).

Although there are various reports on ordered molecular assemblies, including BTB, on solid substrates, the resulting molecular arrangements are highly dependent on substrate features. In order to obtain the desired structures, selection of substrate materials is restricted or surface treatment is necessary, sometimes leading to sacrifice function. Here, we also confirmed that successful transfer of LINAS-1 (Figure S15) and sequential layer-by-layer growth by repeated deposition (Figure 3f and Figure S16) can be made to numerous other substrates besides Si, such as quartz, gold, indium−tin oxide, and graphite. This is a
significant merit of our methodology as the assembled networks can be utilized on desired substrates without suffering from any structural changes. For example, nanosheets obtained by deposition on holey substrates can be used as separation membranes, while those transferred onto electrode substrates can be integrated in thin-film devices such as thin-film transistors and sensors. As shown in Figure 1, the perfectly oriented honeycomb network of LINAS-1 cannot be created on Si substrates by other techniques such as dip casting (Figure 1d), drop casting (Figure 1e), and vacuum deposition (Figure 1f). Similar honeycomb networks of BTB have been observed by scanning tunneling microscopy (STM),

\[ \text{Figure 1e}, \] and vacuum deposition (Figure 1f). Similar honeycomb networks of BTB have been observed by scanning tunneling microscopy (STM),

\[ 30-33 \] but only on specific atomically flat conductive substrates such as graphite, Au(111), and Ag(111). Moreover, the thickness is restricted to that of a monolayer, mainly for the aim to obtain clear images by STM measurements. This contrasts with the fine-tuning of the sheet thickness achieved by our technique. Linear increase of the absorbance with increasing number of deposition cycles (Figure 3f (inset)) provides unambiguous evidence that each deposition step leads to layering of uniform nanosheets with retention of the preferentially oriented structure of multiply deposited LINAS-1, confirmed by GIXRD (Figures S17 and S18). Such layer-by-layer growth allows exquisite tuning of film thickness at the nanoscale.

LINAS-1 is also stable as a free-standing film formed through suspension on a micromesh holey-grid support (Figure 3g). Powdered LINAS-1 can be prepared in large quantities by filtration using a membrane filter—the XRD profile is identical to that of LINAS-1 on Si (Figure S19), again confirming that the porous crystalline network survives even in the absence of any surface support. Finally, LINAS-1 is endowed with exceptional stability to both heating and aging. Both in situ and ex situ XRD measurements and thermogravimetric and differential thermal analysis (TG-DTA) reveal that the perfectly oriented crystalline structure of Si-deposited LINAS-1 is thermally stable to \( \sim 200^\circ \) C despite the nanometer-scale thickness of the sheets (Figures S23 and S24). Notably, an exothermic peak corresponding to the structural change observed by diffraction is evident in the DTA curve of LINAS-1 at \( 230^\circ \) C—this peak is also observed for the BTB bulk crystals, but at a lower temperature (\( 187^\circ \) C). This implies that the structure with perfect molecular orientation in LINAS-1 is more stable than that of the bulk crystal with nonoriented interpenetrated motif. Similarly, GIXRD profiles of LINAS-1 on Si kept in air for three years are identical to those of the same as-prepared sample (Figure S21) establishing its extraordinary stability to long-term aging in the atmosphere.

**Gas Sorption Characteristics.** The LINAS-1 porosity was evaluated by \( \text{N}_2 \)-sorption isotherm measurements at 77 K (Figure 4a). Bulk BTB crystals show an IUPAC type-III shaped isotherm,

\[ \text{Figure 4b}, \] characteristic of a closely packed structure in agreement with the crystal structure determination. On the other hand, the LINAS-1 isotherm exhibits a combined type-I/type-IV shape revealing the presence of both micro/meso- and macropores in LINAS-1. The sharp increase in \( \text{N}_2 \) uptake at the low-pressure region (\( P/P_0 = 0-0.02 \)) implies that the \( \text{N}_2 \) molecules are introduced into the micro/macropores of the nanosheet. The micro/mesopore radii, \( r_p \), extracted from Barrett–Joyner–Halenda (BJH) analysis\n
\[ \text{35} \] (see the Methods) of the \( \text{N}_2 \)-sorption isotherm range between 1 and 3 nm (Figure 4b), in agreement with the effective pore size of \( \sim 2.7 \) nm of the proposed structural model (Figure 2d). The second sharp increase in \( \text{N}_2 \) uptake at high relative pressures (\( P/P_0 = 0.96-1.00 \)) is characteristic of “house-of-cards” inter-nanosheet porosity, commonly present in sheet-shaped materials; the wide distribution in pore radii with a peak maximum at 54 nm in the BJH plot (Figure 4b) reflects the “house-of-cards” porosity.

Complementary water-vapor sorption measurements of LINAS-1 at 298 K reveal the absence of any sharp increase in water uptake (Figure 4c), implying lack of hydrophilic sites in the honeycomb network. The behavior contrasts markedly with that of bulk BTB, which shows a sharp increase in water vapor sorption. Even though the \( \text{N}_2 \)-sorption isotherm of the BTB bulk is consistent with a closely packed structure, the crystal possesses small penetrable cavities along the (101) direction with a diameter of 0.27–0.28 nm (Figure 5f), which exactly match the kinetic diameter of water molecules (2.65 Å) but are smaller than that of \( \text{N}_2 \) (3.64 Å). We consider that free carboxylic groups in the BTB bulk crystal surface attract the water molecules and guide them into the pore cavities. On the other hand, LINAS-1 has neither such effective hydrophilic free carboxylic sites in the open-pore direction nor size-matched pores to water molecules, resulting in featureless water-sorption behavior. This also attests to the high quality and porosity in LINAS-1.

We also examined \( \text{O}_2 \) gas sorption of LINAS-1 at 77 K (Figure 4d). The isotherm in bulk BTB crystals is flat throughout the measured pressure region up to the relative pressure of 0.9. In contrast, we observe a sharp increase in \( \text{O}_2 \) uptake at the low-pressure region in the LINAS-1 isotherm that implies strong interaction between \( \text{O}_2 \) molecules and LINAS-1 frameworks. We consider that in this case nonpolar \( \text{O}_2 \) molecules preferably attach to the hydrophobic pore inner surfaces. This also supports the formation of the high quality nanosheets with nanoporosity (\( \sim 2.7 \) nm) large enough to accommodate \( \text{O}_2 \) molecules (kinetic diameter, 3.46 Å) in contrast to the bulk material in which the cavity diameter is only 0.27–0.28 nm.

**CONCLUSIONS**

The results presented here show the first successful construction of a single-crystal-like noncovalent organic framework nanosheet (named LINAS-1) by facile interferential synthesis using molecular building blocks. In situ synchrotron X-ray diffraction at the air/liquid interface established the immediate formation of highly crystalline LINAS-1 through the spontaneous organization of the molecular components upon their spreading onto the water surface at room temperature. LINAS-1 is completely oriented; the 2D honeycomb networks are parallel to and the \( \pi-\pi \)-stacked molecular columns are vertical to the water surface, while the corresponding bulk crystal adopts a nonoriented multiply interpenetrated structure. Such interpenetration is often seen in large framework materials, in which different sublattices occupy the same space resulting in the minimization of the void volume and the collapse of the porous motif. By virtue of using a liquid interface as an assembly work-platform, both the molecular orientation of the building components and the growth direction are finely controlled, suppressing the formation of intrinsically nonporous complex interpenetrated structures. The existence of both size- and position-regulated permanent effective pores was demonstrated by gas-sorption measurements.

LINAS-1 nanosheets display additional superior quality characteristics. They are exceptionally robust to both heating and environmental aging and are remarkably versatile; they can be transferred to diverse types of substrates such as silicon, quartz, gold, graphite, and holey grids while preserving their highly crystalline perfectly oriented structure intact. The versatility of the interferential synthetic strategy opens the way for designing-in functional features to nanosheets derived from classes of materials.
that are not restricted to those with the potential for exfoliation of their bulk crystalline layered structures.

**METHODS**

Nanosheet Preparation (LINAS-1). A PTFE-coated Langmuir trough was filled with pure water or acidic water as a subphase. The surface of the subphase was carefully cleaned by mild surface-touch vacuuming. A 1 mM solution of BTB in mixed toluene/methanol solvent (2:1, v/v) was spread onto the water subphase with a microsyringe. Surface pressure–area (II–A) isotherm measurements were performed at a continuous barrier pressing speed of 500 μm/s at room temperature. The surface pressure was measured with the Wilhelmy plate method. Rapid increase of the surface pressure under surface compression implies formation of a BTB nanosheet at the liquid/air interface (LINAS-1).

Deposition of LINAS-1 on Solid Substrates. A horizontal dipping method was employed for the deposition of LINAS-1 on solid substrates such as quartz, Si, Au/Cr/Si, HOPG, and microgrid. The substrate with interface (plate method. Rapid increase of the surface pressure under surface compression implies formation of a BTB nanosheet at the liquid/air interface (LINAS-1).

**In Situ Grazing-Incidence (GI) Synchrotron X-ray Diffraction (GIXRD) Measurements at Air/Liquid Interfaces.** In situ GI synchrotron XRD measurements of LINAS-1 at the air/liquid interface were performed at room temperature with the six-circle diffractometer on beamline ID10B (Grenoble, France). The dedicated PTFE Langmuir trough (460 × 170 × 5 mm3) mounted on the diffractometer was equipped with a single movable barrier for film compression. LINAS-1 was prepared as described above. The surface pressure was kept constant during individual GIXRD measurements. The Langmuir trough was mounted on an active antivibration system and was enclosed inside an airtight acrylic case with polystyrene windows. Water-saturated helium gas was introduced into the case. The incidence angle, α, for the GIXRD measurements was set at 0.12°. The scattered X-rays were recorded by a one-dimensional (1D) gas-filled position-sensitive detector with vertically located counting wires (VANTEC). The in-plane XRD profiles were collected by scanning over the in-plane θ angle, and the vertical (out-of-plane) scattered intensity was recorded at each 2θ angle. In order to improve the 2θ resolution in the in-plane direction and reduce background contribution, a Soller collimator (0.08°) was placed in front of the 1D detector. The average sheet domain size was estimated with Scherrer’s equation. For the out-of-plane XRD measurements, the incident angle, α, was also fixed at 0.12° and the scattered X-rays were recorded by scanning over the exit out-of-plane θ angle.

Synchrotron XRD Measurements on Solid Substrates. Synchrotron XRD data of LINAS-1 transferred on Si substrates were collected at room temperature with a multidetector diffractometer at beamline BL13XU of SPring-8. Helium gas was supplied through the cell during the measurements. Each data set was recorded using a scintillation counter. The GIXRD measurements were carried out at θ = 0.155°. Diffraction from the sample surface was observed in the in-plane direction (Figure 3B, background subtracted) with the same detection angle as θ. A Soller collimator (0.4°) was placed in front of the detector. The out-of-plane scans were carried out in θ = 2θ scattering geometry (Figure 3C).

**Atomic Force Microscopy.** AFM images were obtained with a Si cantilever (15 N/m) on an SII SPA400 system at room temperature. Dynamic force mode was applied. The height and lateral diameter analysis was conducted on measurements from more than 140 domains in the AFM images.

**Gas Sorption Isotherm Measurements.** N2, O2, and H2O vapor sorption isotherm measurements of powdered LINAS-1 and of bulk BTB crystals were performed using a BETSORP-max automatic volumetric adsorption system (Bel Japan, Inc.) at 77, 77, and 298 K, respectively. The samples were placed in a glass tube and dried under vacuum at 80 °C for 3 h (powdered LINAS-1) or at 150 °C for 24 h (bulk BTB) for solvent removal. The glass tubes containing the samples were set on the measurement system and heated under high vacuum (below 10⁻³ Pa) at 80 °C for 1 h prior to the measurements. The sample tubes were then introduced into a liquid nitrogen cryostat for the N2 and O2 sorption or a thermostat bath for the H2O vapor sorption measurements. Pressure changes were monitored and converted to N2 or H2O vapor uptake.

The pore radii, r_p, were extracted from Barrett–Joyner–Halenda (BJH) analysis of the N2 adsorption isotherm. The pore distribution curve is obtained by plotting the pore volume change ratio, dV_p/dr_p versus pore radius. The BJH analysis is generally applied for cylinder-shaped mesopore-containing materials.

**ASSOCIATED CONTENT**

**3 Supporting Information**

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsnano.7b04447.

Experimental details; bulk single crystal XRD analysis data; UV-vis and IR absorption spectra, PL spectra, powder XRD data, AFM and SEM images (PDF)

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**Notes**

The authors declare no competing financial interest.

The X-ray crystal structure of BTB is deposited in the Cambridge Crystallographic Data Centre (CCDC 1453166).

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