Critical-state parameters for an unsaturated residual sandy clay

D. G. TOLL* and B. H. ONG†

This paper presents experimental data from constant water content triaxial tests on a residual soil from Singapore, tested under unsaturated conditions with measurements of matric suctions. The data for critical-state conditions from these tests are presented within a framework previously proposed by the first author. The unsaturated critical state requires five parameters: \( M_s \), \( M_b \), \( \Gamma_b \), \( \lambda_a \) and \( \lambda_b \). At high degrees of saturation (in a non-aggregated condition) \( M_s = M_b = M \) (the saturated critical-state stress ratio). At lower degrees of saturation, the value of \( M_b \) drops considerably and reaches values near to zero at degrees of saturation below 60%. Conversely, the value of \( M_s \) appears to increase as the degree of saturation reduces. A similar pattern is observed for \( \lambda_a \) and \( \lambda_b \). The functions relating the critical-state parameters to degree of saturation (or volumetric water content) can be expressed in normalised form by referencing them to the saturated state. The same normalised functions for \( M_s/M_x \), \( M_b/M_x \), \( \lambda_a/\lambda_s \) and \( \lambda_b/\lambda_s \) fit the experimental data for the residual sandy clay from Singapore and a lateritic gravel from Kenya previously tested. This suggests that the form of these functions may be common to a range of soil types.

KEYWORDS: compressibility; fabric/structure of soils; partial saturation; residual soils; shear strength; suction

INTRODUCTION

The critical-state concept has been well established as a useful framework within which saturated soil behaviour can be interpreted (Schofield & Wroth, 1968). The critical state is the state achieved by a soil when it exhibits no changes in stress or volume when it is sheared; this is typically achieved at large strains. A soil that is looser (sometimes referred to as wetter) than the critical state will contract in volume under shear to achieve the critical state, or if volume change is prevented then a decrease in pore water pressure will be generated so that the effective stress state moves towards the critical state. A soil that is denser (or drier) than the critical state will dilate, or if volume change is prevented then a decrease in pore water pressure will be generated.

The behaviour of saturated soils is controlled by effective stresses, and for the saturated case water content and volume are interrelated. Therefore the saturated critical state can be expressed through the deviator stress, \( q \), the mean effective stress, \( p' \), and the specific volume, \( v \). At the critical state these variables are related through three critical-state parameters, \( M, \Gamma \) and \( \lambda \):

\[
q = Mp' \quad (1)
\]
\[
v = \Gamma - \lambda \ln p' \quad (2)
\]

where \( M \) is the slope of the projection of the critical-state line in \( q-p' \) space, \( \Gamma \) is the intercept (at \( p' = 1 \) kPa), and \( \lambda \) is the slope of the projection of the critical-state line in \( v - \ln p' \) space.

However, unsaturated soils have an additional phase (the air phase), and it is therefore no longer possible to interpret their behaviour through effective stresses, nor to assume that water content and volume are linked. For unsaturated soils, the stress state can be represented by two stress state variables, the net stress (\( \sigma - u_w \)) and the matric suction (\( u_w - u_a \)) (Fredlund & Morgenstern, 1977), where \( u_w \) is the pore air pressure and \( u_w \) is the pore water pressure. In addition to specific volume (or void ratio), the phase state of the soil has to be represented by an additional variable: this can be either gravimetric water content (\( w \)), volumetric water content (\( \theta \)) or degree of saturation (\( S_w \)).

Toll (1990) proposed that the critical state for unsaturated soils could be expressed in terms of \( q, p - u_a, u_w - u_a, v \) and \( S_w \). The unsaturated critical state requires five parameters, \( M_w, M_b, \Gamma_{sb}, \lambda_a \) and \( \lambda_w \):

\[
q = M_w(p - u_a) + M_b(u_w - u_a) \quad (3)
\]
\[
v = \Gamma_{sb} - \lambda_a \ln (p - u_a) - \lambda_w \ln (u_w - u_a) \quad (4)
\]

Note that the contributions due to suction have been labelled \( M_b \) and \( \lambda_w \) (rather than the terminology \( M_s \) and \( \lambda_s \) originally used by Toll, 1990). This is for consistency with the use of \( \phi^b \) by Fredlund et al. (1978) to represent the friction angle for changes in matric suction. The terminology for the

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contributions due to net stress ($M_a$ and $\lambda_b$) are already consistent with the use of $\phi^a$ by Fredlund et al. (1978).

Toll (1990) showed that, for a compacted lateritic gravel, $M_a$, $M_b$, $\lambda_a$ and $\lambda_b$ did not have unique values for unsaturated conditions. It was shown that the variation in the parameters could be expressed as functions of the degree of saturation, $S_r$. It was argued that it was not degree of saturation per se that affected these parameters, but rather the fabric of the soil. Clayey materials compacted dry of optimum moisture content develop an aggregated or ‘packet’ fabric, and these effects can be observed even in a predominantly granular material that has less than 10% clay. Toll (2000) showed that, for compacted soils, the degree of saturation could give an indication of the amount of aggregation present.

This paper presents test data from triaxial tests on a residual soil from Singapore, tested under unsaturated conditions with measurements of matric suction. The data for critical-state conditions from these tests are presented within the framework proposed by Toll (1990). The critical-state parameters for the Singapore residual soil are established and are compared with the pattern of behaviour for Kiunyu gravel reported by Toll (1990). Their variation is considered with respect to degree of saturation and also volumetric water content.

**CRITICAL-STATE STRESS RATIOS FOR UNSATURATED SOILS**

Toll (1990) showed that, for a lateritic gravel (Kiunyu gravel), the values of the critical-state stress ratios $M_a$ and $M_b$ varied in a manner that could be expressed as functions of the degree of saturation, $S_r$. Subsequent research has also identified variations in the critical-state parameters due to changes in suction. The results of tests on compacted kaolin by Zakaria et al. (1995) and Wheeler & Sivakumar (1995) showed that the slope of the critical-state line in $q$ vs. $u_s$ space (that is, $M_a$) increased as suction increased. This would be consistent with an increase in $M_a$ with decreasing degree of saturation (as degree of saturation will reduce as suction increases). It is also consistent with other observations that the friction angle for changes in net stress (that is, $\phi^a$) increased with increasing suction (Escario & Saez, 1986), although these latter observations were not made for critical-state conditions. However, Delage et al. (1987) and Maatouk et al. (1995) found that, for silts, $\phi^a$ decreased with increased suction. Ng et al. (2000) found that $M_b$ seemed to be constant and independent of the value of suction for two loosely compacted volcanic fills. Nevertheless, their values of $M_b$ did show a clear variation with degree of saturation, similar to that shown by Toll (1990).

Figure 1 shows the functions for $M_a$ and $M_b$ as they relate to degree of saturation for Kiunyu gravel (Toll, 1990). Toll (2000) also expressed the critical-state parameters for Kiunyu gravel in terms of the friction angles $\phi^a$ and $\phi^b$, and a scale for friction angles is shown in Fig. 1. These represent friction angles at the critical state.

At lower degrees of saturation (dry of optimum water content), $M_a$ has greater values than $M_b$ (the stress ratio in saturated conditions). This was explained by Toll (1990, 2000) as being due to the presence of aggregations at lower degrees of saturation, causing the soil to behave in a coarser fashion than would be justified by the grading. In unsaturated conditions the aggregated fabric can be maintained during shear because the suction gives strength to the aggregations. In a saturated soil, any aggregations that did exist would be broken down during shear and would not be expected to affect the critical-state parameters. In any case, samples compacted at high degrees of saturation (above optimum water content) would not be expected to have an aggregated fabric.

The same mechanism could explain the drop in $\phi^a$ for silts observed by Delage et al. (1987) and Maatouk et al. (1995). Silty materials would tend to desaturate at much lower suctions. At intermediate values of suction the aggregated fabric would be held together by the suction and could give high angles of friction. However, as the suction increases, the aggregations can progressively desaturate. The desaturated parts would no longer be held together by suction, and would break down. In this case the value of $\phi^a$ would drop as the aggregated fabric broke down. This same mechanism is likely to affect clayey soils, but at much higher values of suction.

It can be seen from Fig. 1 that, at low degrees of saturation, $M_b$ becomes significantly lower than $M_a$ and eventually drops to zero. It would be likely that the packet fabric also affects $M_b$ as the water phase withdraws into the packets and suction makes no contribution to the overall strength of the soil. It is also possible that this represents a similar effect to the reduction in the friction angle for suction ($\phi^b$) as suction increases, as observed by Escario & Saez (1986) and Fredlund et al. (1987). This effect can be explained by a reduction in the area of water over which the suction acts (very like the $c$ factor used by Bishop, 1959). It has been shown by Fredlund et al. (1995) and Vanapalli et al. (1996) that the reduction of $\phi^b$ can be related to normalised volumetric water content or degree of saturation.

The variation in the $M_b$ parameter with degree of saturation has been modelled using a function similar to those used by Vanapalli et al. (1996). The function is

$$\frac{M_b}{M_a} = \left(\frac{S_r - S_{r1}}{S_{r2} - S_{r1}}\right)^k$$

The values for $M_b/M_a$ are numerically very similar to the ratio $\tan \phi^b/\tan \phi^a$ (as used by Fredlund et al. (1995) and Vanapalli et al. (1996)), and can be treated as equivalent, so the same function can be used whether stress ratios or friction angles are being used. The function applies only for $S_{r2} < S_r < S_{r1}$; when $S_r > S_{r1}$ the ratio $M_b/M_a$ is taken as 1, and when $S_r < S_{r2}$ it is taken as 0.

Vanapalli et al. used full saturation (that is, 100%) as the first reference state, $S_{r1}$, and the degree of saturation at residual suction as the second reference state, $S_{r2}$. Here, the two reference states have been chosen to fit the experimental
data. To achieve this, values of $S_1 = 88\%$, $S_2 = 50\%$ and $k = 2.0$ have been used, and the resulting function is shown in Fig. 1. $S_1$, is the point where $M_b = M_s$, and $S_2$ represents the point where $M_b = 0$. The significance of these points will be discussed later. The value of $k = 2.0$ provides the required degree of curvature of the function between the two reference states. Fredlund et al. (1995) suggested that $k$ should be 1-0 or greater. A value of 2.0 is similar to that used by Vanapalli et al. (1996), where typically values of 2.2–2.5 were found to be appropriate for a glacial till. 

A similar function has also been used to model the variation in $M_s$. The function used is

$$
\frac{M_a}{M_s} = \left(\frac{M_a}{M_s}_{\text{max}}\right) - \left(\frac{M_a}{M_s}_{\text{max}} - 1\right) \left(\frac{S_1 - S_2}{S_1 - S_2}\right)^k
$$

(6)

This gives a function that varies from $M_a/M_s = 1$ at $S_1$, up to a maximum value of $(M_a/M_s)_{\text{max}}$ at $S_2$. To fit the experimental data, the same values of $S_1 = 88\%$ and $S_2 = 50\%$ were appropriate (the same values used for fitting $M_b$) with $(M_a/M_s)_{\text{max}} = 1.25$. The value of $k = 1.0$ was used, since a linear relationship gives a better fit with the experimental data.

### CRITICAL-STATE COMRESSIBILITIES FOR UNSATURATED SOILS

The slopes of the critical-state surface in specific volume space, $\lambda_a$ and $\lambda_b$, like the stress ratios, do not have unique values. Toll (1990) suggested that they can be related to degree of saturation. He also suggested that $\lambda_a$ in equation (4) could be related to $\Gamma_a$ (the intercept for saturated conditions) through the degree of saturation:

$$
\lambda_a = 1 + \frac{\Gamma_a - 1}{S_1}
$$

(7)

Alonso et al. (1990) suggested that the compressibility in $v$, $p - u_s$ space (in their terminology, $\lambda(s)$) could be related to the value at zero suction ($\lambda(0)$) either through the logarithm of suction or by an exponential function of suction. This would suggest that the compressibility for net stress would reduce as suction increases. Experimental data from Matyas & Radhakrishnan (1968) and Cui & Delage (1996) show this to be the case.

For saturated soils it has been observed that the slope of the critical-state line is parallel to the normal consolidation line. If we assume that this is also true for unsaturated soil, the previous arguments concerning the normal consolidation surface can be applied to the slope of the critical-state line. This would suggest that $\lambda_a$ would reduce with an increase in suction (that is, it would reduce as the degree of saturation reduces). This is not borne out by the observation of Toll (1990) that $\lambda_a$ would increase with a reduction in degree of saturation, nor is it supported by other experimental evidence. Wheeler & Sivakumar (1995) found that the slope of the critical-state line in the net stress plane increased as suction increased. Ng et al. (2000) showed that $\lambda_a$ was considerably steeper for a suction of 80 kPa than for the zero suction condition. It is not clear whether this discrepancy means that the slope of the critical-state surface is different from that of the normal consolidation surface, or whether an increase in $\lambda_a$ with a reduction in degree of saturation (or increase in suction) is due to the fabric of compacted soils.

Figure 2 shows the functions for $\lambda_a$ and $\lambda_b$ as they relate to degree of saturation for Kiunyu gravel (Toll, 1990). It was argued that the presence of aggregations (and hence a more open fabric) causes the soil to behave in a more compressible fashion (with respect to net stress) than is justified by the grading. Therefore at lower degrees of saturation (greater aggregation) $\lambda_a$ had greater values than $\lambda_s$ (the compressibility in saturated, non-aggregated conditions). At low degrees of saturation $\lambda_a$ becomes significantly lower than $\lambda_s$, and at low water contents drops to zero. This could be due to the packet fabric, reflecting the fact that the water phase withdraws into the packets and suction makes no contribution to the compressibility of the soil. It is also possible that the effect is due to degree of saturation per se, reflecting the reduced normalised area of water over which the suction acts, making the suction less effective in causing volume change.

### TEST PROGRAMME

Constant water content triaxial tests were carried out on unsaturated samples of the residual soil from Singapore. In addition, consolidated drained tests on saturated samples were performed to define the saturated parameters. The soil tested was obtained from samples taken at a site adjacent to the Nanyang Technological University, School of Civil and Structural Engineering (NTU-CSE), which is located in the Jurong sedimentary deposit of Singapore. Samples were taken from a depth of 4–5 m. Prior to testing, the soil was air-dried for two weeks and passed through a sieve with an aperture size of 2:36 mm. Some classification properties of the soil (after removal of material < 2:36 mm) are given in Table 1, and the grain size analysis is given in Fig. 3. Comparable information is also given for Kiunyu gravel, tested by Toll (1990).

The Jurong soil was mixed at a water content of 15-6% (1-4% wet of optimum water content) and statically compacted into a mould to obtain specimens of 100 mm height and 50 mm diameter. The intention was to produce specimens with a bulk density of 1-97 Mg/m$^3$ and a dry density of 1-71 Mg/m$^3$ (void ratio of 0-6). The compaction was...

### Table 1. Classification properties

<table>
<thead>
<tr>
<th></th>
<th>Jurong residual soil</th>
<th>Kiunyu gravel</th>
</tr>
</thead>
<tbody>
<tr>
<td>Liquid limit: %</td>
<td>36</td>
<td>56–66</td>
</tr>
<tr>
<td>Plastic limit: %</td>
<td>22</td>
<td>27–31</td>
</tr>
<tr>
<td>Plasticity index: %</td>
<td>15</td>
<td>29–35</td>
</tr>
<tr>
<td>Specific gravity of solids: %</td>
<td>2.72</td>
<td>3.2</td>
</tr>
</tbody>
</table>

Fig. 2. Variation of critical-state compressibilities with degree of saturation for Kiunyu gravel [$\nu = \Gamma_{ab} - \lambda_a \ln (p - u_s) - \lambda_s \ln (u_s - u_a)$]
carried out in 10 layers, with each layer having an equal thickness of 10 mm. Full details of the specimen preparation are given by Ong (1999).

Three drained triaxial tests were performed on saturated samples of the Jurong soil. The samples were consolidated isotropically to effective confining stresses of 50, 100 and 150 kPa. The shearing rate used was 0.009% min, the same as that used in the constant water content tests on unsaturated samples. The results are plotted in Fig. 4.

A series of 15 constant water content triaxial tests were carried out on the Jurong soil. The axis-translation technique was used to control and measure the matric suction. In this technique an elevated pore air pressure is used, so that the pore water pressure becomes positive. A high air entry porous disc is used so that the pore water pressure can be measured independently of the pore air pressure. In this test series, a porous ceramic with an air entry value of 500 kPa was used. Pore water pressure was measured at the base of the specimen. A diffused air volume indicator as described by Fredlund & Rahardjo (1993) was used to remove air bubbles from the pore water pressure measuring system before shearing.

After compaction, specimens were initially saturated by applying an elevated back pressure. The specimens were then allowed to equilibrate under appropriate values of isotropic confining pressure, $\sigma_3$, pore air pressure, $u_a$, and pore water pressure, $u_w$, as shown in Table 2. Net confining stresses ($\sigma_3 - u_a$) were set at 50, 150 or 250 kPa. Matric suctions ($u_a - u_w$) varied from 170 to 400 kPa. The specimens were then sheared under constant water content conditions (that is, preventing any flow of water in or out of the specimen). However, the specimens could still change in volume owing to flow of air. The specimen volume was determined by measuring the volume of flow in or out of

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Fig. 3. Particle size distributions for Jurong soil and Kiunyu gravel

Fig. 4. Results for drained tests carried out on saturated samples, plotted against axial strain: (a) deviator stress, $q$; (b) volume strain
the Perspex triaxial cell surrounding the specimen using a
GDS pressure/volume controller. A correction was applied to
allow for the downward movement of the loading ram.

A constant strain rate of 0.009% min\(^{-1}\) was applied to the
specimens. This rate was chosen after a pilot study by Ong
(1999) showed that this rate (and lower) gave consistent
readings of matric suctions in constant water content tests.

The results for tests carried out at a net stress of 50 kPa
are plotted in Fig. 5. Similar plots for tests carried out at
a net stress of 150 kPa are shown in Fig. 6. The results
for tests carried out at a net stress of 250 kPa are shown in
Fig. 7.

**THE SATURATED CRITICAL STATE FOR JURONG SOIL**

The saturated critical state provides a useful reference for
the unsaturated critical state. Therefore, initially data from
saturated tests and those for constant water content tests that
achieved high degrees of saturation (greater than 90% at the
end of the tests) will be considered in order to give indications
of the saturated critical-state values. At high degrees of
saturation the pore air phase exists as occluded air bubbles,
and it is acceptable to analyse the soil behaviour from
unsaturated tests in terms of effective stresses. This provides
additional information to extend the findings obtained from
the limited number of tests carried out on saturated samples.

The five constant water content tests that reached high
degrees of saturation at the end of the tests were cw50-170,
cw50-200, cw150-220, cw150-240 and cw250-280. These
tests achieved degrees of saturation between 91% and 93%.
Therefore the results of these five tests can be analysed in
terms of effective stresses. The results for cw50-170 and
 cw50-200 (with net stress of 50 kPa and suctions of 170 and
200 kPa respectively) are shown in Fig. 5. Results for
cw150-220 and cw150-240 (with net stress of 150 kPa and
suctions of 220 and 240 kPa respectively) are shown in Fig.
6, and the result for cw250-280 (with net stress of 250 kPa
and suction of 280 kPa) is shown in Fig. 7.

The saturated drained tests (cd50, cd100 and cd150) (Fig.
4) show that both the deviator stress and the volume strain
levelled off at axial strains of 12–15%, meaning that a true
critical state is achieved. It can be seen that for the constant
water content tests, cw50-170 and cw50-200 (Fig. 5), the
volume strain shows a tendency to level off at an axial strain
above 10%, suggesting that the critical state is being
approached. However, the deviator stress and the pore water
pressure are still changing towards the end of the tests. A
similar picture is seen for cw150-220 (Fig. 6) in cw250-
280 (Fig. 7).

**Table 2. Initial stress values for the constant water content tests**

<table>
<thead>
<tr>
<th>Test</th>
<th>(\sigma_3)</th>
<th>(u_a)</th>
<th>(u_w)</th>
<th>(\sigma_3 - u_a)</th>
<th>(u_w - u_a)</th>
</tr>
</thead>
<tbody>
<tr>
<td>cw50-170</td>
<td>500</td>
<td>450</td>
<td>280</td>
<td>50</td>
<td>170</td>
</tr>
<tr>
<td>cw50-200</td>
<td>500</td>
<td>450</td>
<td>250</td>
<td>50</td>
<td>200</td>
</tr>
<tr>
<td>cw50-230 (1)</td>
<td>500</td>
<td>450</td>
<td>220</td>
<td>50</td>
<td>230</td>
</tr>
<tr>
<td>cw50-230 (2)</td>
<td>500</td>
<td>450</td>
<td>220</td>
<td>50</td>
<td>230</td>
</tr>
<tr>
<td>cw50-250</td>
<td>500</td>
<td>450</td>
<td>150</td>
<td>50</td>
<td>100</td>
</tr>
<tr>
<td>cw50-300</td>
<td>500</td>
<td>450</td>
<td>50</td>
<td>50</td>
<td>0</td>
</tr>
<tr>
<td>cw150-220</td>
<td>600</td>
<td>450</td>
<td>230</td>
<td>150</td>
<td>220</td>
</tr>
<tr>
<td>cw150-240</td>
<td>600</td>
<td>450</td>
<td>210</td>
<td>150</td>
<td>240</td>
</tr>
<tr>
<td>cw150-260</td>
<td>600</td>
<td>450</td>
<td>190</td>
<td>150</td>
<td>260</td>
</tr>
<tr>
<td>cw150-300</td>
<td>600</td>
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<td>150</td>
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<td>300</td>
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<tr>
<td>cw150-350</td>
<td>600</td>
<td>450</td>
<td>100</td>
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<td>350</td>
</tr>
<tr>
<td>cw150-400</td>
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<td>450</td>
<td>50</td>
<td>150</td>
<td>400</td>
</tr>
<tr>
<td>cw250-280</td>
<td>700</td>
<td>450</td>
<td>170</td>
<td>250</td>
<td>280</td>
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<tr>
<td>cw250-400</td>
<td>700</td>
<td>450</td>
<td>50</td>
<td>250</td>
<td>400</td>
</tr>
</tbody>
</table>

**Fig. 5. Results for constant water content tests at net confining stress of 50 kPa and various matric suctions, plotted against axial strain:**
(a) deviator stress, \(q\); (b) change in pore water pressure, \(\Delta u\); (c) volume strain.
A true critical state (that is, no change in deviator stress, pore pressure or volume) was rarely achieved by the end of the constant water content tests. However, the rate of change of these parameters did generally reduce considerably towards the end of the test. In some cases the deviator stress continued to increase (and the pore water pressure to decrease), but by the end of the test the stress path followed a constant stress ratio ($q/p^9$). This can be seen in the stress paths shown in Fig. 8. It is always difficult to define a true critical-state condition for bonded or dense materials as they often fail through the development of distinct shear surfaces, where non-homogeneous deformations (and the possibility for clay particle alignment) will affect the results. Nevertheless, for the purpose of this discussion it will be assumed that the ultimate stress ratio ($q/p^9$) defines the critical-state condition.
stress ratio, \( M \). The saturated tests (cd50, cd100, cd150) clearly define \( M \), as they failed in a ductile fashion.

The \( q/p' \) ratios for the three saturated drained tests and five constant water content tests at high degrees of saturation are shown in Fig. 9. The saturated drained tests show a consistent ultimate \( q/p' \) ratio with a value of 1.23. There is considerable variation in the ultimate stress ratios for the five constant water content tests, as can be seen in Fig. 8. The values range from 0.94 in test cw150-240 to 1.37 in test cw50-170. There does seem to be a trend of reducing \( q/p' \) ratio with increasing suction, but with these limited data it is difficult to be sure. It is likely that higher suctions produce more brittle behaviour (and hence the formation of distinct failure surfaces), and this is reflected in the lower \( q/p' \) values. However, the ultimate value of \( q/p' \) for the saturated drained tests of 1.23 will be taken as the saturated value of \( M \).

The volumetric behaviour of the three saturated drained tests and the five constant water content tests at high degrees of saturation is shown in Fig. 10. The end point of each test is shown by a ◊ symbol. A line is drawn on the figure representing the projection of the critical-state line (CSL) in \( e - \log p' \) space. The line drawn represents values of \( \Gamma = 2.0 \) and \( \lambda = 0.08 \).

The three saturated drained tests (cd50, cd100 and cd150) show contractive behaviour throughout. Similarly, the four constant water content tests (that start above the CSL (cw50-170, cw50-200, cw150-220 and cw250-280) show contractive behaviour. Test cw150-240, which starts with a void ratio denser than the critical state, shows strong dilative behaviour in the latter stages of the test. This difference in state for test cw150-240 (that is, starting below the CSL) explains why this test shows such a different stress path in Fig. 8.

In conclusion, the test results for saturated tests and those at high degrees of saturation provide values for the saturated critical-state parameters of \( M = 1.23 \), \( \Gamma = 2.0 \) and \( \lambda = 0.08 \).

THE UNSATURATED CRITICAL STATE FOR JURONG SOIL

Table 3 shows the values of the stress state variables \( (q, p - u_k \text{ and } u_k - u_w) \) and the phase state variables \( (w, e \text{ and } S) \) at the end of each test. It can be seen that the degree of saturation at the end of the tests varies from 93% down to 61%.

By examining the values of \( q, p - u_k \text{ and } u_k - u_w \), it is possible to assess values of \( M_e \) and \( M_b \) from equation (3). It was assumed that these values would relate to the degree of saturation, as suggested by Toll (1990). An initial assumption was made that \( M_e \) was constant and equal to \( M_a \), independent of \( S \). Values of \( M_b \) were then determined so that the predicted value of deviator stress, \( q \), calculated from equation (3) (based on the observed values of \( p - u_k \) and \( u_k - u_w \)) was equal to the observed value of \( q \). This initial assessment gave a distribution of values of \( M_b \) that did not show a consistent function with degree of saturation. The values of \( M_e \) were then adjusted by trial and error and new values of \( M_b \) calculated so that reasonably consistent functions of both \( M_a \) and \( M_b \) were achieved. The results of this assessment are shown in Fig. 11.

Although there is some scatter in Fig. 11, the predominant trends are reasonably clear. The values of \( M_a \) and \( M_b \) for the tests at high degrees of saturation are seen to be equal to each other, and range from 0.94 to 1.37, scattering around the value for \( M_e \) of 1.23 as discussed in the previous section. As the degree of saturation reduces, the values of \( M_a \) drop considerably, and reach values less than 0.1 at a degree of saturation near 60%. Conversely, the values of \( M_b \) appear to increase with reducing degree of saturation. This is consistent with the pattern observed by Toll (1990).

The functions shown by lines in Fig. 11 are given by equation (5) using the same values of \( S_{t1} = 88\%, \ S_{t2} = 50\% \) and \( k = 2.0 \) as used for Kiunyu gravel, but are now based on an \( M_e \) value of 1.23 appropriate for the Jurong soil. It can be seen that the functions give a very reasonable fit to the data obtained for the Jurong sedimentary soil. This suggests that the form of the function may be common to a range of soil types, as the same form fits both a lateritic gravel and a residual sandy clay. These functions, expressed as the normalised variables \( M_a/M_e \) and \( M_b/M_e \), are plotted in Fig. 12.

By examining the values of \( e, p - u_k \text{ and } u_k - u_w \), it is possible to assess values of \( \lambda_a \) and \( \lambda_b \) from equation (4). Values of \( \Gamma_{ab} \) were calculated from the value of \( \Gamma_a \) and \( S_e \) using equation (7). As before, it was assumed that the values of \( \lambda_a \) and \( \lambda_b \) would relate to the degree of saturation. The same process used to determine \( M_e \) and \( M_b \) was used to...
determine the functions for $\lambda_a$ and $\lambda_b$. The results of the assessment are given in Fig. 13. The functions shown by lines in Fig. 13 are the same functions as observed for Kiunyu gravel (Fig. 2) but replotted for $\lambda_s$ value of 0.08. It can be seen that the functions correspond very closely to the data obtained for the Jurong sedimentary soil. Like the functions for $M_a = M_s$ and $M_b = M_s$, this suggests that the form of the functions for $\lambda_a = \lambda_s$ and $\lambda_b = \lambda_s$ may be common to a range of soil types. These functions, expressed as the normalised variables $\lambda_a = \lambda_s$ and $\lambda_b = \lambda_s$, are plotted in Fig. 14.

FUNCTIONS USING VOLUMETRIC WATER CONTENT

The functions for the critical-state stress ratios, expressed as the normalised variables $M_a/M_s$ and $M_b/M_s$ (Fig. 12), have so far been expressed in terms of degree of saturation using equations (5) and (6). The form of these functions has also been investigated in terms of volumetric water content. This has the advantage that it can be related to the soil–water characteristic curve when expressed in terms of volumetric water content, an approach that has been favoured by

<table>
<thead>
<tr>
<th>Test</th>
<th>$q$</th>
<th>$p - u_a$</th>
<th>$u_a - u_w$</th>
<th>$w$</th>
<th>$e$</th>
<th>$S_r$</th>
</tr>
</thead>
<tbody>
<tr>
<td>cw50-170</td>
<td>311</td>
<td>154</td>
<td>73</td>
<td>19.4</td>
<td>0.581</td>
<td>90.5</td>
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<td>cw50-200</td>
<td>283</td>
<td>144</td>
<td>85</td>
<td>19.2</td>
<td>0.569</td>
<td>91.7</td>
</tr>
<tr>
<td>cw50-230(1)</td>
<td>175</td>
<td>108</td>
<td>158</td>
<td>16.7</td>
<td>0.686</td>
<td>66.3</td>
</tr>
<tr>
<td>cw50-230(2)</td>
<td>247</td>
<td>132</td>
<td>158</td>
<td>16.7</td>
<td>0.637</td>
<td>71.3</td>
</tr>
<tr>
<td>cw50-250</td>
<td>332</td>
<td>161</td>
<td>182</td>
<td>16.2</td>
<td>0.692</td>
<td>63.8</td>
</tr>
<tr>
<td>cw50-300</td>
<td>204</td>
<td>118</td>
<td>186</td>
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<td>0.628</td>
<td>63.1</td>
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<tr>
<td>cw50-400</td>
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<td>150</td>
<td>292</td>
<td>14.1</td>
<td>0.528</td>
<td>72.8</td>
</tr>
<tr>
<td>cw150-220</td>
<td>385</td>
<td>278</td>
<td>44</td>
<td>18.7</td>
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<td>91.0</td>
</tr>
<tr>
<td>cw150-240</td>
<td>433</td>
<td>294</td>
<td>168</td>
<td>17.0</td>
<td>0.505</td>
<td>91.6</td>
</tr>
<tr>
<td>cw150-260</td>
<td>463</td>
<td>304</td>
<td>171</td>
<td>16.5</td>
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<tr>
<td>cw150-300</td>
<td>391</td>
<td>280</td>
<td>178</td>
<td>15.9</td>
<td>0.583</td>
<td>74.3</td>
</tr>
<tr>
<td>cw150-350</td>
<td>480</td>
<td>310</td>
<td>282</td>
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<td>0.481</td>
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<tr>
<td>cw150-400</td>
<td>579</td>
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<td>cw250-280</td>
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<td>45</td>
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<tr>
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<td>483</td>
<td>241</td>
<td>14.5</td>
<td>0.649</td>
<td>60.9</td>
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</tbody>
</table>
Fredlund et al. (1995) and Vanapalli et al. (1996). They suggested that

$$\frac{\tan \phi^b}{\tan \phi^a} = \left( \frac{\theta - \theta_s}{\theta_a - \theta_s} \right)^k$$

(8)

A similar approach will be used here, represented by the function:

$$\frac{M_b}{M_s} = \left( \frac{\theta - \theta_{p2}}{\theta_{p1} - \theta_{p2}} \right)^k$$

(9)

The values of $\theta_{p1}$ and $\theta_{p2}$ are chosen as the best fit for the experimental data. Their significance will be discussed in the following.

The $M_s$ values were modelled using the following function (similar to equation (6)):

$$\frac{M_a}{M_s} = \left( \frac{M_a}{M_s \text{ max}} \right) - \left( \frac{M_a}{M_s \text{ max}} - 1 \right) \Theta^k$$

(9)

The values of $\theta_{p1}$ and $\theta_{p2}$ that best fitted the Kiunyu gravel data (for both $M_b$ and $M_s$) were $\theta_{p1} = 0.27$ and $\theta_{p2} = 0.43$. For the Jurong soil, $\theta_{p1} = 0.18$ and $\theta_{p2} = 0.33$. In both cases a value of $k = 2.0$ was used for $M_b$ and $k = 1.0$ for $M_s$ (the same values as used for fitting degree of saturation data). The results are plotted in terms of volumetric water content in Fig. 15, and can be seen to give a realistic fit to the data. The data for both materials are also plotted in terms of normalised water content, $\Theta$, in Fig. 16. The curvature of the $M_b$ function can be clearly seen (modelled with $k = 2.0$) compared with the linear function for $M_s$ (using $k = 1.0$). It can be seen that both sets of data fit the same functions in terms of normalised volumetric water content, as did the degree of saturation functions (Fig. 12).

Soil–water characteristic curves determined from pressure plate tests are shown for both Kiunyu gravel and Jurong soil in Fig. 17. The test equipment in both cases was limited to a maximum matric suction of 500 kPa. This means that the residual suction cannot be defined as it is clear that both soil–water characteristic curves continue to show significant changes in volumetric water content up to the limit of the test data. However, the curve for the Jurong soil does show the slope of the curve decreasing at suctions above 200 kPa.

The values of $\theta_{p1}$ and $\theta_{p2}$ for the Jurong soil have been indicated in Fig. 17. It can be seen that $\theta_{p2}$ is close to the air entry value for the soil, possibly representing the end of the boundary effect zone and the start of the primary transition zone as defined by Vanapalli et al. (1996). $\theta_{p2}$ coincides with the break in slope of the curve at 200 kPa, possibly representing the end of the primary transition zone. Unfortunately, the limit on suction for the test on Kiunyu gravel means that only the initial part of the curve is identified. The value of $\theta_{p1}$ is known and falls on the initial part of the soil–water characteristic, again possibly representing the start of the primary transition zone. The value of $\theta_{p2}$ falls outside the experimental results, so no conclusions can be drawn for the Kiunyu gravel.

It is clear that the two reference states ($\theta_{p1}$ and $\theta_{p2}$) needed to define the function for variations in the stress ratios do not coincide with saturation ($\theta_s$) and residual conditions ($\theta_i$), as suggested by Fredlund et al. (1995) and Vanapalli et al. (1996). From the limited information on soil–water characteristic curves it would seem that the limits
of volumetric water content appropriate for normalisation could be defined by the start and end of the primary transition zone.

CONCLUSIONS

The critical state for unsaturated soils can be expressed in terms of the following variables; deviator stress, \( q \); net mean stress, \( p \); specific volume, \( v \); and degree of saturation, \( S \). The unsaturated critical state requires five parameters: \( M_a \), \( M_b \), \( \Gamma_{ab} \), \( \lambda_a \) and \( \lambda_b \). These can be referenced to the saturated critical-state parameters \( M_a' \), \( M_b' \) and \( \lambda_a' \). The variation in critical-state stress ratios (\( M_a' \) and \( M_b' \)) and the critical-state compressibilities (\( \lambda_a' \) and \( \lambda_b' \)) can be expressed as functions of the degree of saturation of the soil.

Results from constant water content tests on unsaturated samples of soils from the Jurong residual soil in Singapore have been used to determine the critical-state parameters for the soil in an unsaturated state. Tests on saturated specimens indicate saturated critical-state parameters of \( M_a' = 1.23 \), \( \Gamma_{ab} = 2.0 \) and \( \lambda_a' = 0.08 \). At high degrees of saturation \( M_b' \) drops considerably, and reaches a value near to zero at degrees of saturation below 60%. Conversely, the value of \( M_a' \) appears to increase with reducing degree of saturation. This is consistent with the pattern observed by Toll (1990).

A similar pattern is observed for \( \lambda_a' \) and \( \lambda_b' \). At high degrees of saturation \( \lambda_a' = \lambda_b' = \lambda_a \). As the degree of saturation reduces the value of \( \lambda_b' \) drops, and reaches a value near to zero at degrees of saturation below 60%. Conversely, the value of \( \lambda_a' \) clearly increases significantly with reducing degree of saturation, as observed by Toll (1990).

The functions relating the critical-state parameters to degree of saturation can be expressed in normalised form by referencing them to the saturated state. The same normalised functions for \( M_a/M_a' \), \( M_b/M_b' \), \( \lambda_a/\lambda_a' \) and \( \lambda_b/\lambda_b' \) fit the experimental data for the Jurong residual soil from Singapore and also the data for a lateritic gravel from Kenya reported by Toll (1990). This suggests that the form of these functions may be common to a range of different soil types.

The functions can also be expressed in terms of normalised volumetric water content. It seems from limited data that the two reference states needed for normalisation may be defined by the start and end of the primary transition zone identified from the soil–water characteristic curve.

REFERENCES


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NOTATION

\( e \) void ratio
\( k \) fitting parameter used for relating shear strength to \( \Theta \)
(Fredlund et al., 1995)
\( M_a, M_b \) critical-state stress ratio for saturated conditions (= \( q/p' \))
\( M_a, M_b \) critical-state stress ratio for changes in net mean stress
\( p \) mean stress (= \( \sigma_1 + \sigma_2 + \sigma_3 \)/3)
\( q \) deviator stress (= \( \sigma_1 - \sigma_3 \))
\( S \) degree of saturation
\( u_a \) pore air pressure
\( u_w \) pore water pressure
\( \theta \) change in pore water pressure
\( v \) specific volume (= 1 + \( e \))
\( w \) gravimetric water content
\( \Gamma, \Gamma_b \) critical-state intercept for volumetric plane at \( p' = 1 \) kPa
\( \Gamma_{ab} \) critical-state intercept for volumetric planes at
\( p - u_a = 1 \) kPa and \( u_a - u_w = 1 \) kPa
\( \theta_v \) volumetric water content
\( \theta_v \) volumetric water content at saturation
\( \theta_{rs} \) volumetric water content at residual suction
\( \theta_{rp} \) volumetric water content at start of primary transition zone
\( \theta_{rs} \) volumetric water content at end of primary transition zone
\( \lambda, \lambda_b \) slope of critical-state line in \( v - \ln p' \) plane for saturated conditions
\( \lambda_a \) slope of critical-state line in \( v - \ln (\sigma - u_a) \) plane
\( \phi \) angle of friction for changes in net stress (\( \sigma - u_a \))
\( \phi_b \) angle of friction for changes in matric suction (\( u_a - u_w \))

Fig. 17. Soil–water characteristic curves for Kiunyu gravel and Jurong soil


