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# Forming a Single Molecule by Magnetoassociation in an Optical Tweezer

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We demonstrate the formation of a single NaCs molecule in an optical tweezer by magnetoassociation through an *s*-wave Feshbach resonance at 864.11(5) G. Starting from single atoms cooled to their motional ground states, we achieve conversion efficiencies of 47(1)%, and measure a molecular lifetime of 4.7(7) ms. By construction, the single molecules are predominantly [77(5)%] in the center-of-mass motional ground state of the tweezer. Furthermore, we produce a single *p*-wave molecule near 807 G by first preparing one of the atoms with one quantum of motional excitation. Our creation of a single weakly bound molecule in a designated internal state in the motional ground state of an optical tweezer is a crucial step towards coherent control of single molecules in optical tweezer arrays.

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Ultracold polar molecules, with their tunable long-range interactions and rich internal structures, provide a promising means for quantum simulation of novel phases of matter [1–5] and quantum information processing [6–10]. Many key ingredients of these proposals, such as the dipolar exchange interaction [11], long coherence times of nuclear spin and rotational states [12–14], and information transduction between different molecular degrees of freedom [15], have been demonstrated utilizing molecular gases and ions. To realize the aforementioned applications, coherent control of individual ultracold molecules is needed, at the level of single quantum states in both the internal and motional degrees of freedom. A new generation of molecular experiments thus aims towards systems that are simultaneously scalable and able to provide a high level of control over individual particles. This is being pursued through molecular ions in ion traps [16,17] and through neutral molecules in optical tweezers, both directly cooled [18] and assembled from their constituent laser-cooled atoms [19–23].

In the bottom-up approach of molecular assembly, forming a single weakly bound molecule is an important milestone towards creating arrays of rovibrational ground-state molecules with tunable interactions. Previously, a single weakly bound molecule (NaCs  $a^3\Sigma, v = -1$ ) was created from an atom pair by two-photon Raman transfer, but suffered from rapid photon scattering that subsequently scrambled its internal state [21]. Here, we form a weakly bound molecule by magnetoassociation through a Fano-Feshbach resonance (FR), which has been established as a robust technique to bridge bi-alkali atoms to rovibrational ground-state molecules, including in bulk gases [24–29]

and in optical lattices [30–33]. While creating molecules in optical tweezers would provide the additional benefits of flexibility and configurability, the peak intensity of optical tweezers is in general several orders of magnitude higher than in optical lattices at the same trapping frequencies; this presents a potential obstacle to the adiabatic magnetoassociation of atoms into molecules.

In this Letter, we demonstrate that a single Feshbach molecule can be formed in an optical tweezer by magnetoassociation of individually trapped atoms. The molecule has a lifetime of a few milliseconds, limited by scattering from the trap light. By controlling the motional states of the atoms, we can control both the motional and the rotational state of the molecule produced.

The experiment begins with a single  $^{23}\text{Na}$  atom and a single  $^{133}\text{Cs}$  atom loaded stochastically from a dual-species magneto-optical trap into separate optical tweezers. The atoms are imaged after loading so that we can postselect on whether both species are initially loaded (two-body) or only one is loaded (one-body). The atoms are then simultaneously cooled to their respective 3D motional ground states by polarization gradient cooling and Raman sideband cooling. Details of the trapping and cooling procedures have been reported previously [21,34]. After the atoms are cooled to their motional ground states, we prepare them in the lowest Zeeman energy level: Na  $|F = 1, m_F = 1\rangle$  and Cs  $|F = 3, m_F = 3\rangle$ . This choice of hyperfine channel eliminates the possibility of spin-changing inelastic collisions, which could occur in our previous work [35], and allows production of stable Feshbach molecules.

Our search for Na-Cs FRs was initially guided by multichannel quantum defect theory, using singlet and

triplet scattering lengths previously determined from interaction shift spectroscopy [35]. In the present work, we have developed a coupled-channel model of Na + Cs. Singlet and triplet potential curves were obtained by adjusting the interaction potentials in Ref. [36] to reproduce the binding energy of the least-bound triplet state [35] and the position of the  $s$ -wave resonance described below. The adjustment procedures were similar to those used for K + Cs in Ref. [37]. The resulting potential curves were combined with the spin-spin dipolar interaction and an estimate of the 2nd order spin-orbit coupling to allow calculations on states with nonzero relative angular momentum. Scattering calculations were carried out using the MOLSCAT package [38,39] and bound-state calculations using the BOUND and FIELD packages [39,40].

In experiments with bulk gases, FRs are usually detected through enhanced three-body loss [41]. The tweezers, however, contain only two atoms, in their lowest-energy states; this precludes three-body loss and spin-changing inelastic collisions. We therefore directly utilize molecule formation by magnetoassociation for the FR search [42]. The experimental sequence for magnetoassociation is shown schematically in Fig. 1(a). After single atom trapping, ground-state cooling and hyperfine state preparation, a magnetic field produced by a pair of Helmholtz coils is ramped up in 40 ms along the axial direction of the optical tweezers to 866.5 G. The two traps are then merged, so that the Na and Cs atoms are held in a single optical tweezer at 1064 nm and a peak intensity of 81 kW/cm<sup>2</sup>, giving trapping frequencies  $\omega_{\text{Cs}} = 2\pi \times (30, 30, 5)$  kHz and  $\omega_{\text{Na}} \approx 1.07\omega_{\text{Cs}}$  [21]. The magnetic field is then ramped linearly down to various values at a rate of 1 G/ms. If the magnetic field ramp crosses a FR then magnetoassociation is possible. For detection, the tweezer is separated back into the species-specific tweezers before ramping the magnetic field down to zero for imaging the surviving atoms, as shown by the solid line in Fig. 1(a). Because the imaging detects only the atoms, magnetoassociation to form Feshbach molecules is manifest as a two-body loss.

We locate an  $s$ -wave Feshbach resonance at 864.11(5) G, as shown in the lower panel of Fig. 1(b); the position is determined by a fit to an error function. An additional loss feature is detected at 864.5 G, which we attribute to photoassociation enhanced by a narrow resonance nearby [43]. As confirmation of the two-body nature of the processes, we also measure the survival rates of the single atoms when loaded without the presence of the other species; these are shown in the upper panel of Fig. 1(b), and show no features. The contrast between the left and right asymptotes in the two-body loss data gives a molecule conversion efficiency of 47(1)%.

The atom-to-molecule conversion process can be described by a Landau-Zener (LZ) type avoided crossing with an efficiency that depends on the ramp rate of the

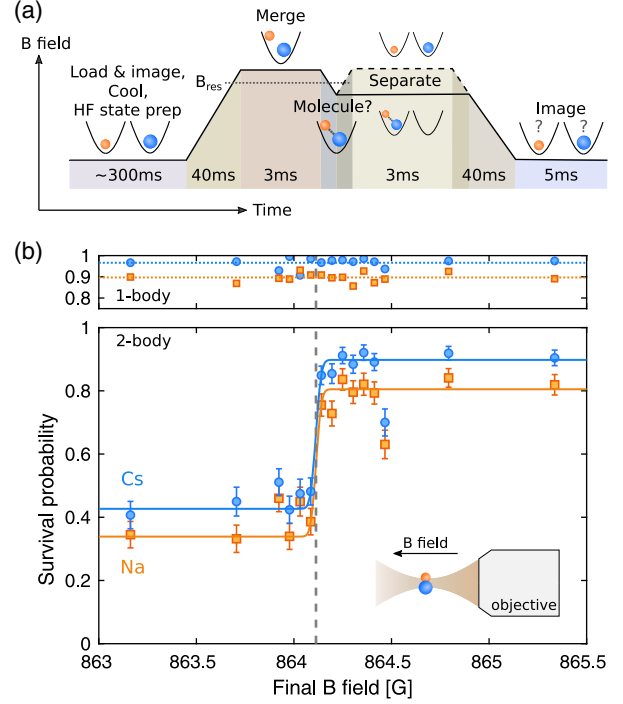


FIG. 1. Magnetoassociation at the  $s$ -wave resonance. (a) Schematic of magnetic field ramp and trap merge and separation sequence as a function of time. Solid (dashed) line indicates one-way molecule conversion (conversion back to atoms for molecule survival detection). Time spent for magnetoassociation is varied for different experiments; see text for details. (b) Determination of the  $s$ -wave resonance location. The magnetic field is ramped linearly from 866.5 G to the various magnetic fields at 1 G/ms. Lower panel: survival probabilities of Na (orange squares) and Cs (blue circles) when both species are loaded. The solid lines are fits to an error function, from which we extract the left and right asymptote values and resonance location. Vertical dashed line indicates resonance location determined from the fit. Upper panel: same experimental run with initial one-body loading. Horizontal dotted lines are the mean values for each species.

magnetic field and characteristic parameters intrinsic to the FR [31,42,44]. To investigate molecule formation, we vary the rate of a linear magnetic field ramp from 866 to 863.9 G. The resulting joint Na and Cs survival probabilities are shown as the purple circles in Fig. 2(a). A lower two-body survival probability indicates a higher molecule conversion probability. The one-way molecule conversion efficiency follows the LZ formula  $p_{\text{mol}} = 1 - e^{-2\pi\delta_{\text{LZ}}}$ , where  $\delta_{\text{LZ}} = (2\pi n_2/\mu)|a_{\text{bg}}\Delta/\dot{B}|$  [41]. Here  $\Delta = 1.29$  G and  $a_{\text{bg}} = 30.7 a_0$  are the width and background scattering length of the Feshbach resonance, obtained from coupled-channel calculations using the method of Ref. [45],  $\mu = 19.60$  amu is the reduced mass,  $n_2 = \int \int n_{\text{Na}}(\mathbf{r})n_{\text{Cs}}(\mathbf{r})d\mathbf{r}$  is the density of a single pair of Na and Cs atoms in the optical tweezer, and  $\dot{B}$  is the magnetic field ramp rate, which is varied experimentally. The purple curve in Fig. 2(a) is the best fit to the LZ formula. The fit value of the pair density

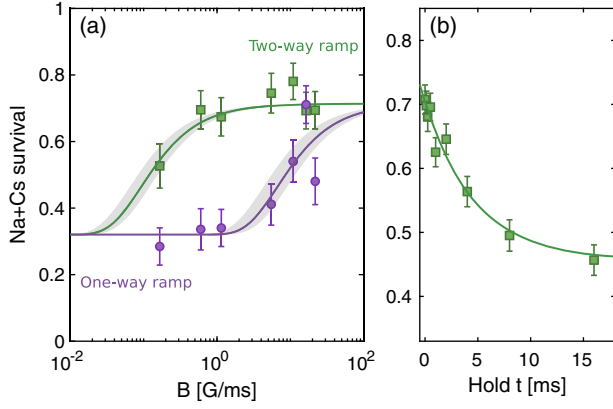


FIG. 2. Molecule formation and dissociation efficiencies. (a) Purple circles indicate Na + Cs joint survival probability after magnetoassociation. The magnetic field is ramped linearly from 866 to 863.9 G at different rates. Green squares indicate Na + Cs joint survival probability with an additional reverse magnetic field ramp at the same rate after molecule formation. The solid lines are best-fit curves and the gray shaded areas indicate the errors on the fit. See text for fit details. (b) Lifetime of Feshbach molecule held at  $B - B_{\text{res}} = -0.3$  G and trap intensity 81 kW/cm<sup>2</sup>. Solid line is best fit to an exponential decay.

$n_2 = 2.5(9) \times 10^{13}$  cm<sup>-3</sup> is in good agreement with our trap parameters.

To detect the survival of the Feshbach molecules in the optical tweezer, we dissociate the molecules back into atoms by performing a reverse magnetic field ramp, as shown by the dashed line in Fig. 1(a). We assume the Feshbach molecule dissociates with certainty since no molecular state exists above resonance [46]. The two-way conversion efficiency back to atoms is limited by the time the Feshbach molecules spend in the optical tweezer. This can be expressed as  $p_{\text{atom}} \approx e^{-t_{\text{mol}}/\bar{\tau}}$ , where  $t_{\text{mol}} = (2|B - B_{\text{res}}|/\dot{B})$  is the time spent below the FR, and  $\bar{\tau}$  is the molecular lifetime, averaged over the ramped magnetic field. We can directly measure the lifetime of the molecules at a particular magnetic field in a separate experiment by holding the molecules for varying times before dissociating and detecting atom survival. At  $B - B_{\text{res}} = -0.3$  G and a trap peak intensity of 81 kW/cm<sup>2</sup>, we observe a lifetime of  $\tau = 4.7(7)$  ms as shown in Fig. 2(b). The green curve in Fig. 2(a) is a best fit to the two-way ramp that yields  $\tau = 6(2)$  ms, which agrees well with the lifetime measurement and corroborates the lifetime-limited conversion efficiency. This lifetime is promising for future work such as coherently transferring to the rovibrational ground state by stimulated Raman adiabatic passage, which would take tens of microseconds [25].

In order to characterize the factors limiting the lifetime of the Feshbach molecules in the optical tweezers, we measure the lifetimes under various hold conditions. In one case, we vary the power of the trap used to hold the molecules after formation. The Feshbach molecules are

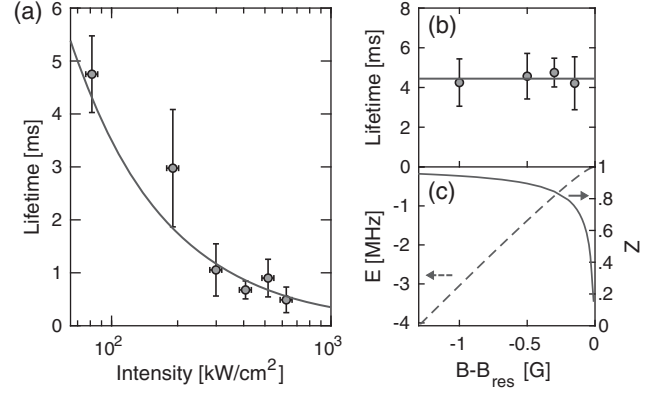


FIG. 3. Characterization of *s*-wave Feshbach molecule lifetime. (a) Dependence on trap intensity. The trap is adiabatically ramped to and held at different intensities after magnetoassociation at  $B - B_{\text{res}} = -0.3$  G. The line is a best fit to inverse scaling. (b) Dependence on magnetic field. The field is ramped to different values during magnetoassociation while the trap intensity is fixed at 81 kW/cm<sup>2</sup>. The solid line is the mean value. (c) The negative binding energy  $-E_b$  (dashed line, left axis) and closed-channel fraction  $Z$  (solid line, right axis) from coupled-channel bound-state calculations as a function of magnetic field. See text for details.

formed and dissociated with a ramp rate of 3 G/ms and are held at  $B - B_{\text{res}} = -0.3$  G. We find that the lifetime of the molecules is inversely proportional to trap intensity, as shown in Fig. 3(a), suggesting that the lifetime is limited by scattering from the trap light. This observation agrees with that previously reported for Feshbach molecules in optical lattices [31]. From the scattering rates, we determine the imaginary part of the polarizability at 1064 nm to be 2.8(3) Hz/(kW/cm<sup>2</sup>); this is  $\sim 100$  times higher than expected from theoretical predictions [47]. A similar excessive scattering is also observed in excited NaCs molecular states. There is not yet any clear theoretical explanation of these observations.

We also vary the magnetic field at which the Feshbach molecules are held, over fields that correspond to binding energies up to  $E_b = 3$  MHz. The trap peak intensity is fixed at 81 kW/cm<sup>2</sup>. As shown in Fig. 3(b), we observe no significant variation of the lifetime in this range. The scattering rate of the trapping light depends on the Franck-Condon overlap between the Feshbach molecular state and excited molecular states in the vicinity of the tweezer wavelength; under some circumstances this is proportional to the closed-channel fraction  $Z(B)$  of the wave function for the Feshbach molecule [31]. We have performed coupled-channel bound-state calculations to evaluate  $Z(B)$  from the expression  $Z(B) = (\mu_b - \mu_a)/(\mu_{\text{bare}} - \mu_a)$ , where  $\mu_b$  ( $\mu_{\text{bare}}$ ) is the magnetic moment of the Feshbach molecular state (the bare molecular state well below threshold), and  $\mu_a$  is that of the separated atoms [41].  $Z$  and  $-E_b$  are shown as functions of magnetic field in Fig. 3(c). From these we find that this resonance has only a small region of universality.



At the magnetic fields we use,  $B - B_{\text{res}}$  between  $-1$  and  $-0.15$  G,  $Z$  is close to 1 and varies slowly with magnetic field.

The conversion efficiency of an atom pair to a single Feshbach molecule and the motional state of the resulting molecule are both determined by the motional state of the atom pair, described in terms of the relative and center-of-mass (c.m.) motions [48]. Atom pairs can be most efficiently converted to molecules when they are in the ground state of relative motion. We directly measure the population of the relative ground state by interaction shift spectroscopy in a separate experiment, as described in Ref. [35]. We find a relative motional ground-state probability of  $\sim 58\%$ , which combined with hyperfine-state preparation fidelities (Na  $\sim 88\%$ , Cs  $\sim 96\%$ ) is consistent with our molecule conversion efficiency of  $45\%$ – $50\%$ , depending on experimental conditions [48].

The c.m. motional state of the Feshbach molecule is inherited from that of the constituent atoms. An atom pair that is in its relative motional ground state but an excited c.m. motional state may still be magnetoassociated to form a molecule. We can infer the atomic c.m. ground state population from independent Raman sideband thermometry measurements of each atom, and estimate that  $77(5)\%$  of the resulting molecules are in the c.m. motional ground state [48]. It should be noted that both the molecule conversion efficiency and the c.m. ground state population of the Feshbach molecules are not fundamentally limited to their present level, and can be increased by improved atomic ground-state cooling fidelity.

In addition to controlling the motional state of the Feshbach molecule, we can control the internal state through choice of the atomic motional states. In particular, we use a  $p$ -wave resonance to form a rotationally excited molecule. While atoms in their motional ground states have no relative angular momentum, we can controllably excite the radial motional state of the Na atom by one motional quantum so that the relative motional state of the pair is  $\sim 24\%$  in the excited state [48]. Since the excitation is in the plane perpendicular to the magnetic field axis, as shown in the inset of Fig. 1(b), the resulting state has relative angular momentum  $M_L = \pm 1$ .

Our coupled-channel calculations predict two  $p$ -wave bound states that cross threshold near 807 G, with total molecular spin angular momentum  $M_{F,b} = 4$  and 5. Each of these splits into components with total angular momentum  $M_{\text{tot}} = M_{F,b}$  and  $M_{F,b} \pm 1$ . The colliding atoms have  $m_{F,\text{Na}} + m_{F,\text{Cs}} = 4$  and  $M_{\text{tot}} = 3$  or 5 in the radially excited motional state. We thus expect 3 resonant features for such atoms. We detect these features by FR-enhanced photoassociation [49], as shown in Fig. 4. The two atoms are held for 20 ms in a tweezer with peak intensity  $1350 \text{ kW/cm}^2$  after merging the traps at a magnetic field value that is scanned. We detect simultaneous two-body loss when the atoms are photoassociated via the excited electronic states

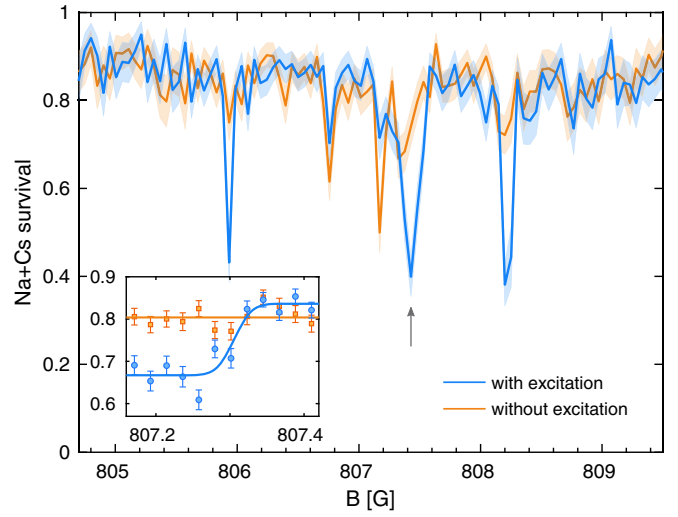


FIG. 4.  $p$ -wave FR spectroscopy by FR-enhanced photoassociation. The atoms are held at a fixed magnetic field in an intense tweezer. Blue (orange) line shows spectrum with (without) motional excitation of Na. Shaded areas indicate error bars. The arrow indicates the resonance used for magnetoassociation in the inset. Inset:  $p$ -wave Feshbach molecule formation. The magnetic field is ramped linearly from 807.6 G to the various magnetic fields at  $0.02 \text{ G/ms}$ . Same color scheme as main figure. Blue (orange) curve is fit to error function (mean value across the range).

by the tweezer light. For comparison, we also show the same scan without the motional excitation on Na.

As for the  $s$ -wave Feshbach molecules, we ramp the magnetic field across a  $p$ -wave Feshbach resonance to transfer the atoms into a  $p$ -wave molecule. The inset of Fig. 4 shows the survival probability when the magnetic field is ramped linearly down from 807.6 G to various fields at a rate of  $0.02 \text{ G/ms}$  in a tweezer held at  $81 \text{ kW/cm}^2$  peak intensity. We observe a clear two-body loss feature when we perform the motional excitation, in contrast to the case of no motional excitation. We attribute this to  $p$ -wave molecule formation and find a conversion efficiency of  $16(2)\%$ .

In conclusion, we have formed single NaCs Feshbach molecules by magnetoassociation in an optical tweezer, using newly identified FRs. In particular, we have formed  $s$ -wave Feshbach molecules in their motional ground state starting from atoms cooled to their motional ground state and  $p$ -wave molecules from atoms prepared in specific excited motional states. Feshbach molecules are not susceptible to three-body collisional losses in optical tweezers, as they are in bulk gases, allowing us to achieve high conversion efficiencies that are not fundamentally limited. While the lifetimes of the Feshbach molecules are limited by scattering from the tweezer light, this does not pose an obstacle to further transfer to the rovibrational ground state by stimulated adiabatic Raman passage. With in-situ atomic rearrangement possible in tweezer arrays [50,51], the

molecule conversion efficiency we have achieved would scale up to a lattice filling fraction near 50%; this is higher than previously achieved [32,33], and would allow studies of percolating many-body dynamics in 2D geometries [52]. Identification of ground-state molecules for trap rearrangement might also be possible by imaging the atoms that are not converted to molecules. The high conversion efficiency and exquisite control over the molecules demonstrated here, combined with the configurability provided by the tweezers, render optical tweezer arrays of dipolar molecules a promising platform for quantum simulation and quantum information processing.

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